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A comparison of time and frequency domain-based approaches to laser Doppler vibrometer instrument vibration correction --Manuscript Draft--

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Abstract:	<p>When laser Doppler vibrometers are used in the presence of ambient vibration, it is essential to compensate for the additional vibration signal content. In practice, compensation is realised by independently determining the instrument vibration and subtracting it from the erroneous measurement. When these vibrations are transient in nature, time domain-based processing must be used to carry out the correction. However, recent implementation of such an approach on stationary signals showed a factor of eight increase in performance over the previously established frequency domain-based alternative. Therefore, the work described in this paper initially focuses on determining the cause of the inconsistency and proposes a revised frequency domain approach. This revised approach offers near-equivalent performance to its time domain-based equivalent, with the latter approach offering only a factor of 0.26 increase in performance. However, despite the advantages of selecting the time domain-based technique, it typically requires high oversampling factors to allow for the accurate synchronisation of the various transducer type signals. Up until now, the only method available to determine the relationship between the sampling frequency and the performance would be experimentally, which is laborious and time consuming. Therefore, the significance of this paper is the development and experimental validation of an analytical model which predicts the sampling frequency dependence of the time domain correction technique performance. Using this, a framework was developed which allows for the optimal implementation of either correction technique and specifies the required acquisition parameters.</p>

1. Revised frequency domain based LDV base vibration correction technique presented.
2. Enhanced signal processing shown to yield factor of seven performance improvement.
3. Recently proposed time domain processing method performance rigorously examined.
4. Model developed to relate correction signal synchronisation error and performance.
5. Model validated against experimental data with excellent agreement.
6. Framework defined for selection of correction technique and hardware specification.

Hello,

We would like to express our gratitude for accepting the paper. The references have been amended, however, if any issues remain please let us know. Additionally, in this manuscript, the credit author statement was removed and now is only in the Word file attached. I hope this isn't an issue. Thank you.

Kind regards,

The authors

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9 A comparison of time and frequency domain-based
10 approaches to laser Doppler vibrometer instrument
11 vibration correction
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28
29 **Abstract**
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32 vibration, it is essential to compensate for the additional vibration signal
33 content. In practice, compensation is realised by independently determining
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Keywords: mobile laser Doppler vibrometry, vibration measurement, non-stationary instrument vibration correction, time domain signal processing, transient vibration

1. Introduction

Laser Doppler vibrometers (LDVs) have become indispensable and widely adopted vibration measurement tools [1]. Increasingly they have been applied to mobile applications, which include buried landmine detection [2-5], terrestrial seismology [6], orbital seismology [7-9], and vibratory health assessment from drones [10]. The integration of LDVs into autonomous vehicles, such as unmanned aerial vehicles (UAVs), is a recent addition to the extensive and ever-growing plethora of LDV applications which has been receiving increased attention in recent years [11, 12]. Such solutions unlock enormous potential for truly autonomous and remote measurement campaigns within traditionally inaccessible, remote or hazardous environments. There has, however, been a particular recent fundamental advance that has enabled the pursuit of this ambitious application domain.

Specifically, this advance relates to correction of the measured signal, where the contribution of the LDV's own vibration - otherwise indistinguishable from the intended measurement - can be completely removed, thereby fully recovering the target surface vibration velocity. Establishing solutions for LDV measurement correction has been completed for: single beam devices in the presence of arbitrary, six degree-of-freedom vibration [13]; scenarios in which beam steering optics, which might vibrate independently of the sensor head, are used [14]; and more recently, for scanning LDVs, where the laser beam scan angle must also be accounted for [15].

Dual beam LDVs can be configured to optically subtract the motion of

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26 the sensor head by mixing the measurement beam from the target with a
10 reference beam from a static surface. However, this correction technique has
11 two main practical limitations in the context of mobile LDV applications.
12 Previous work has shown that to obtain full six degree-of-freedom correc-
13 tion the correction measurement must occur along the beam axis [11]. This
14 would require the reference beam to be colinear and pointing in the opposite
15 direction to the measurement beam and focused on the static reference sur-
16 face. While this might be realisable in the laboratory, for mobile, field-based
17 applications, this is practically impossible to achieve due to the simultane-
18 ous positioning and focusing requirements. Conversely, accelerometers are
19 robust, readily available and enable direct measurement of sensor head vi-
20 bration. Recently developed solutions have, therefore, been focused on their
21 use, thereby offering accessible options for the practising vibration engineer.
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28 In general, independent vibration measurements of the additional velocity
29 contributions to the LDV signal are required and these are obtained using
30 specifically positioned accelerometers. These additional velocity contribu-
31 tions can be due to a combination of sensor head, scanning head and steer-
32 ing optic vibration; for the sake of brevity, all solutions will be referred to as
33 LDV measurement correction. Depending upon the geometry of the partic-
34 ular set-up, components of these measurements are used to compensate for
35 the additional velocity in the direction of the laser beam. Post-processing
36 has been conducted in the frequency domain where the necessary integration
37 of the accelerometer signals for velocity and their synchronisation with the
38 LDV signal is conveniently implemented. Furthermore, frequency domain
39 representation is well-understood and is commonplace within typical vibra-
40 tion measurement and structural dynamic characterisation campaigns where
41 signals are *stationary* in nature.
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47 In more real-world relevant, field-based vibration measurement scenarios,
48 including those involving the integration of LDVs with autonomous vehicles
49 [11, 12], it can be reasonably expected, however, that both the target *and*
50 the instrument vibration signals will be *transient* in nature. For this rea-
51 son, an alternative approach based entirely in the time domain, was recently
52 proposed and its performance for stationary signals compared against the
53 established frequency domain-based equivalent [16]. While both techniques
54 offer a significant improvement in the corrected LDV signal, the time domain-
55 based technique outperformed its counterpart by a factor of approximately
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eight. The performance difference outcome was unexpected, especially considering that previous work has shown frequency domain integration to be the most accurate [17]. This paper will, therefore, explore the reasons for this performance gap before proposing a revised frequency domain-based approach with significantly improved performance.

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In either domain, the quality of the measurement correction is sensitive to signal synchronisation since any error therein will adversely affect the quality of the corrected velocity estimate. Therefore, the development of a model which relates synchronisation error to the velocity estimate error is paramount. Firstly, the time delay estimate will always have an associated uncertainty, however small; this is minimised using a rigorous relative calibration procedure. Secondly, in the time domain, the implementation of the synchronisation is constrained to integer multiples of the time step, therefore, even a perfect delay estimate is unlikely to lead to perfect synchronisation. While interpolation could be used to upsample time domain data to enable sub-time step alignment, this is not always desirable. Therefore, a thorough investigation into the relationship between the time step and the synchronisation error is required such that an optimal sampling frequency can be selected, maximising the performance of the time domain-based technique.

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The model is validated using significantly oversampled experimental data, downsampled to simulate acquisition over a range of sampling frequencies. This method of experimentally obtaining the sampling frequency dependence also enables the comparison of the time and frequency domain-based techniques across an extended range. Specifically, comparing the relative performances leads to the definition of distinct regions, within each of which the measurement correction outcome can be optimised by selecting the appropriate technique. These findings are then generalised, based on two parameters, to determine the minimum sampling frequency necessary for the time domain-based technique to outperform both others. This in turn enables the user to define the optimal hardware characteristics required for a given measurement campaign. This can be crucial when integrating such sensor solutions into autonomous vehicles as payload versus performance is in a delicate balance.

2. Overview of the correction measurement setup

Practically, the correction of LDV measurements in the presence of sensor head vibration involves the use of properly positioned accelerometers to obtain the correction measurements. The number and positioning of these *correction accelerometers* is determined by the specific nature of the optical setup. For a single-beam LDV, a single accelerometer mounted to the rear of the sensor head colinear with the beam axis is required [11]. This relatively simple setup is convenient for the development of new processing techniques. However, the techniques developed here could be easily expanded to more complex setups which require multiple correction accelerometers, such as a scanning LDV [15].

Fig. 1 illustrates the experimental setup used in this work. It is common with that used in previous work [11, 16] and allows for the independent control of both the target and the LDV vibration. Here, the target vibration is the measurement of interest, while the base vibration simulates the effects of instrument vibration on the LDV measurement. Both the target and the base vibrations were realised using electrodynamic shakers independently driven using uncorrelated broadband white noise up to 200 Hz, generated by a Siemens Digital Industries Software Simcenter SCADAS Mobile data acquisition system and accordingly amplified. The base shaker was a Tira Vibration exciter S 51120 amplified by a Tira Vibration BAA 500 and the target shaker was a Brüel & Kjær V201 M4-CE amplified by a Brüel & Kjær LDS LPA100. While a flat shaker/amplifier response over the frequency range of interest may be desirable, it is not essential since the correction algorithm is completely effective, irrespective of level and phase, across the entire frequency range.

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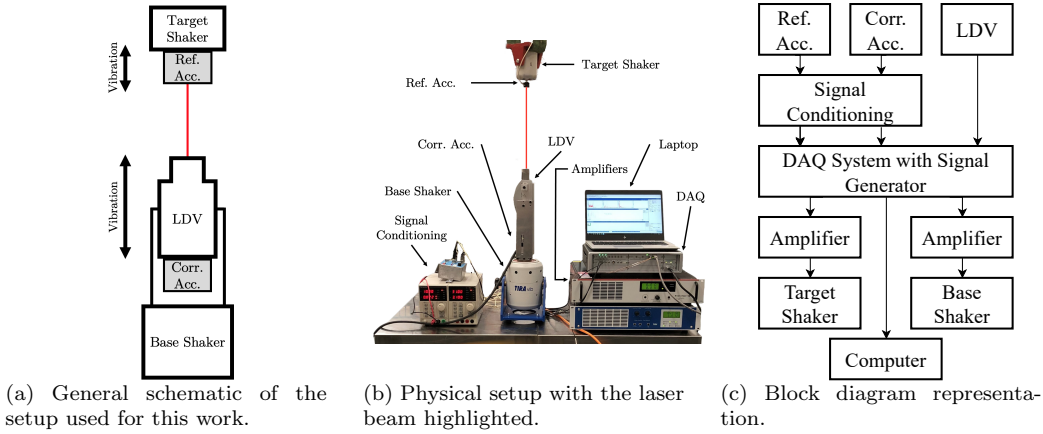


Figure 1: Experimental setup used to simulate a LDV target vibration measurement during base motion vibration. The labels “Corr. Acc.” and “Ref. Acc.” represent the correction and reference accelerometers, respectively.

127 A custom-made aluminium mounting bracket was used to fix a Polytec
 128 NLV-2500-5 Compact Laser Vibrometer to the base motion shaker so that
 129 the laser beam axis was aligned with that of the vibration. An Endevco 770F-
 130 10-U-120 (200 mV/g nominal) DC-response accelerometer was mounted to
 131 the bracket with its sensitive axis colinear with that of the LDV. The tar-
 132 get shaker was suspended directly above the LDV from an overhead crane,
 133 providing isolation from the large base motion shaker. A second Endevco
 134 accelerometer of the same model was mounted to the spigot providing the
 135 ‘true’ vibration measurement. As in earlier work [11], a second, *fixed* LDV
 136 could equally be used for the true vibration measurement. However, one or
 137 both beams would need to be off-axis to enable optical access and this may
 138 require the angular misalignment to be determined and accounted for. In
 139 this work, the use of the reference accelerometer was therefore preferred.

140
 141 There are some practical limitations with the use of accelerometers, in-
 142 cluding the flatness of their amplitude and phase response. Unlike LDVs,
 143 accelerometer performance is typically limited by the first mechanical res-
 144 onance of the mass-spring system. However, they are relatively low cost,
 145 are readily available and can offer acceptable performance in the context of
 146 mobile LDV measurement campaigns which are the focus of the solutions
 147 developed in this body of work. In general, a frequency range from several
 148 Hz to several hundred Hz is considered appropriate with vibration levels on
 149 the order of several tenths to several tens of mm/s. Over such a relatively
 150 narrow frequency range, it is appropriate to compensate for the amplitude

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10 and phase response with a straightforward relative calibration (to the LDV)
11 and this will be described in detail subsequently.
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14 Another limitation of accelerometers is that they typically exhibit a small
15 amount of transverse sensitivity which might degrade correction performance
16 in the presence of significant off-axis vibration, in this case it is only 3% [18].
17 In the experimental setup used here, inevitable rocking motion of the shakers
18 is minimised by centring the mass distribution on the shaker axis and this
19 effect is therefore considered to be negligible. Nevertheless, this and some
20 misalignment between the shaker axes also results in some motion of the
21 LDV beam on the target. While this motion was insufficient to cause the
22 laser beam to deviate substantially from the region of interest on the target,
23 pseudo-vibrations in the LDV signal, which include speckle noise, are asso-
24 ciated with such relative motion of the laser beam across the target surface
25 and these cannot be corrected by the means proposed in this paper. How-
26 ever, combined LDV sensitivity to transverse vibration as a result of both
27 phenomena is on the order of 0.1% [19] and is therefore also considered to be
28 negligible in the context of sensitivity to sensor head vibration [11].
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35 **3. An improved frequency domain-based processing approach**

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37 Accepting the LDV measurement as the reference, the accelerometer sen-
38 sitivities must be adjusted and the signals synchronised prior to being used
39 in post-processing. Both of these require that the accelerometer signals are
40 integrated, however, the integration of a discretised signal commonly leads
41 to the introduction of errors which can manifest themselves as drift. Drift is
42 more readily identified and relatively easily removed in the time domain by
43 subtracting a first order least squares fit. It is, however, less noticeable in
44 the frequency domain and is practically difficult to remove. Conversely, in-
45 tegration is readily implemented and more accurate in the frequency domain
46 [17]. An improved approach to both the relative calibration procedure and
47 the LDV measurement correction is achieved by implementing a combination
48 of time and frequency domain processing techniques.
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53 *3.1. Accelerometer signal relative sensitivity and time delay*

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55 Relative sensitivity determination and time delay estimation first require
56 that a vibration measurement is obtained from all transducers with their
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186 sensitive axes aligned. Practically, this was achieved using an arrangement
187 with the LDV positioned directly above and focused on an accelerometer
188 stack in turn mounted to the spigot of a shaker. Care was taken to eliminate
189 contamination from ambient vibration by placing the entire arrangement on
190 an anti-vibration base.

191
192 The vibration signals are measured and processed according to the pro-
193 cedure shown in Fig. 2 for a single accelerometer channel. In earlier work
194 [11, 13, 15], signals were directly captured as frequency spectra, calculated
195 from Hann-windowed time blocks because the excitation was broadband
196 white noise. Following frequency domain integration of the accelerometer
197 signals, Sensitivity Prefactor and Temporal Alignment values were deter-
198 mined as per the ultimate step in the diagram. In the revised approach, time
199 data are instead acquired with the accelerometer signal immediately con-
200 verted to frequency domain representation, albeit without prior application
201 of a Hann window to the time data blocks. While perhaps considered un-
202 conventional, the lack of windowing is an essential part of the technique as it
203 enables preservation of the overall time domain waveform, thereby allowing
204 more accurate subsequent detrending.

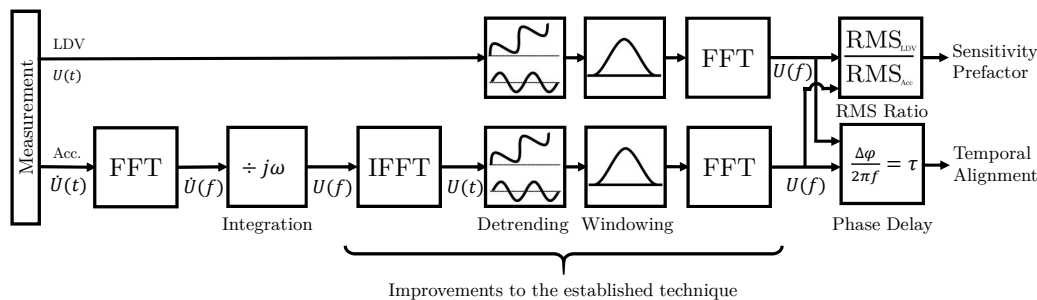


Figure 2: A schematic of the new frequency domain-based relative calibration procedure. The improvements are highlighted by the curly brackets and include the addition of the IFFT, detrending and FFT stages, along with moving the Windowing stage from after the measurement block to just before the second FFT. The signal “Acc.” represents that obtained from either the correction or reference accelerometer.

205 As can be seen in Fig. 2, a $j\omega$ division is used in the frequency do-
206 main to integrate the accelerometer signal. Removal of the resulting drift is
207 achieved by the subtraction of a first order least squares fit from the time
208 domain integrated signal. Since this detrending step will act to remove not

209 only the spurious but also some genuine signal content, the same operation
 210 must be applied to the measured LDV signal. Both signals are now converted
 211 to the frequency domain in the usual way, and implementing a Hann win-
 212 dow on the time data blocks if required. The required Sensitivity Prefactor
 213 and Temporal Alignment parameters are obtained by taking the ratio of the
 214 Root Mean Square (RMS) values and from the phase difference between the
 215 signals, respectively.

216
 217 Fig. 3a shows phase difference plots generated from a single time data
 218 block using the established and the improved frequency domain-based method.
 219 By comparing the two curves, it becomes obvious that detrending leads to
 220 increased agreement between the two types of transducers. As can be seen
 221 in Fig. 3b, the improvement occurs mainly occurs at the lower frequencies.
 222 For a system with a constant phase delay, the group delay can be written as:

$$\Delta\phi = -2\pi f\tau_{\text{meas}} \quad (1)$$

224 where $\Delta\phi$ is the phase difference, τ_{meas} is the measured time delay and f
 225 is the frequency. Therefore, a least squares fit can be used to extract τ_{meas}
 226 from the detrended data set. For this dataset, a value of $\tau_{\text{meas}} = -133.3 \pm 1.8$
 227 μs was obtained and is consistent with equivalent values previously observed
 228 for such sensors and signal conditioning.

229 3.2. Instrument vibration correction

230 *Correction* of the LDV measurement similarly requires integration of the
 231 correction accelerometer signals with detrending therefore being essential for
 232 optimal performance. A revised post-processing approach is shown in Fig. 4.
 233 As for the relative sensitivity adjustment and time delay calculation pro-
 234 cess, the differences between this improved and the previously established
 235 approach are largely captured in the steps to the right of the IFFT and to
 236 the left of the second FFT. Again, the signals are now captured in the time
 237 domain whereas previously frequency spectra were captured directly. In this
 238 case, the integration-related steps are conducted on both the correction and
 239 target reference accelerometers, these having had their relative sensitivities
 240 adjusted and signal time delays estimated. The latter of the two accelero-
 241 meters is only intended for use in the laboratory research campaign where which
 242 provides a ‘true’ vibration measurement for correction performance; for sub-
 243 sequent real-world campaigns, there is no such device since, otherwise, there

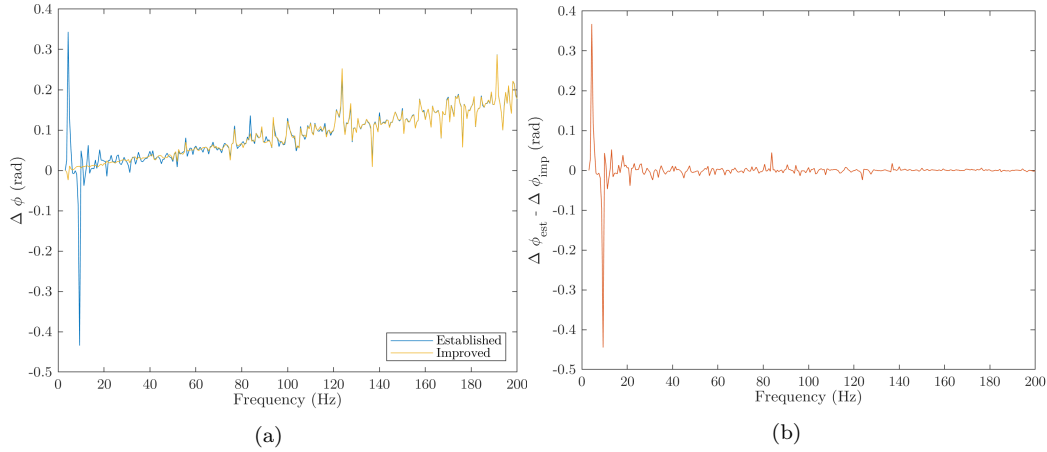


Figure 3: Phase differences for a single, 1.6 s data length, using the established [11], $\Delta\phi_{\text{est}}$, and improved, $\Delta\phi_{\text{imp}}$, frequency domain-based methods; a) phase differences and b) comparison between differences.

244 would be no need to develop the LDV capability for this purpose.

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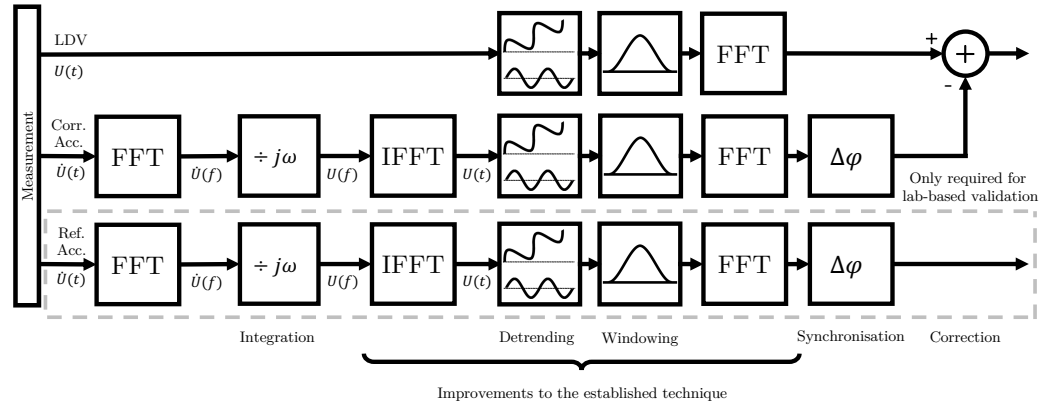


Figure 4: A functional diagram representing the improved frequency domain-based technique. The improvements are highlighted by the curly brackets and include the addition of the IFFT, detrending and FFT stages, along with moving the Windowing stage from after the measurement block to just before the second FFT. The signal “Corr. Acc.” and “Ref. Acc.” represent that of the correction and reference accelerometers, respectively.

246 It is important to also note here that the LDV measurement itself must
 247 be subject to the same detrending step. Otherwise, the corrected LDV signal

248 contains some signal content that has been removed from the detrended ac-
 249 celerometer signals. Following the second FFT, the correction processing is
 250 similar to that previously described [11], the exception being that, here, the
 251 accelerometer signals are already in velocity. Incorporating the previously
 252 determined time delays, before subtracting the correction accelerometer sig-
 253 nal from the LDV (in complex representation), yields the *corrected* LDV
 254 signal for direct comparison with the ‘true’ vibration, given by the reference
 255 channel. The correction performance can be quantified using the error re-
 256 duction, given by [15]:

$$R = -10 \log_{10} \left(\frac{\text{MSE}_{\text{corr}}}{\text{MSE}_{\text{meas}}} \right) \text{ dB} \quad (2)$$

258 where MSE_{meas} and MSE_{corr} are the mean square error of the LDV signal
 259 before and after correction, respectively, when taking the processed reference
 260 accelerometer signal as the ‘true’ vibration signal.

261
 262 Previously, MSE_{meas} and MSE_{corr} have been derived assuming there is no
 263 DC offset in the signal [15]. However, a complete description of the MSE for
 264 a signal of N spectral lines and for the m th spectra would be:

$$\begin{aligned} \text{MSE}_m^{\text{signal}} = & (a_{0,m}^{\text{signal}} - a_{0,m}^{\text{true}})^2 + \\ & \frac{1}{2} \sum_{n=1}^N (A_{n,m}^{\text{signal}} - A_{n,m}^{\text{true}})^2 + (B_{n,m}^{\text{signal}} - B_{n,m}^{\text{true}})^2 \end{aligned} \quad (3)$$

265 where $A_{n,m}^{\text{signal}}$ and $B_{n,m}^{\text{signal}}$ are the real and imaginary parts, respectively,
 266 of either the measured or corrected LDV signal at the n th spectral line for
 267 the m th spectra. The same notation applies to $A_{n,m}^{\text{true}}$ and $B_{n,m}^{\text{true}}$, which are the
 268 reference accelerometer equivalents. Similarly, $a_{0,m}^{\text{signal}}$ and $a_{0,m}^{\text{true}}$ are the DC
 269 component equivalents. When averaging across multiple spectra, the error
 270 reduction takes the following form:

$$R = -10 \log_{10}(\bar{r}) \text{ dB} \quad (4)$$

271 with:

$$\bar{r} = \frac{1}{M} \sum_{m=1}^M \frac{\text{MSE}_m^{\text{corr}}}{\text{MSE}_m^{\text{meas}}} \quad (5)$$

272 where M is the total number of acquired spectra.

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273 *3.3. Measurement correction performance comparison*

274 It is useful to visualise the performance difference between the two tech-
275 niques as a function of the frequency. To do this, the error reduction can
276 be calculated for each spectral line and plotted. The MSE in Eq. (2) can be
277 substituted for the square error, SE, to preserve the frequency information.
278 This plot is improved if a mean of each spectral line, n , is taken across the
279 multiple spectra, m . Algebraically, this is given by:

$$\begin{aligned} \text{SE}_{\text{signal}}(n) &= \frac{1}{M} \sum_{m=1}^M (a_{0,m}^{\text{signal}} - a_{0,m}^{\text{true}})^2 && \text{for } n = 0 \\ &= \frac{1}{2M} \sum_{m=1}^M (A_{n,m}^{\text{signal}} - A_{n,m}^{\text{true}})^2 + (B_{n,m}^{\text{signal}} - B_{n,m}^{\text{true}})^2 && \text{for } n > 0 \end{aligned} \quad (6)$$

280 where all symbols are as previously defined. Substituting the SE in place
281 of the MSE in Eq. (2) would then give:

$$R(n) = -10 \log_{10} \left(\frac{\text{SE}_{\text{corr}}}{\text{SE}_{\text{meas}}} \right) \text{ dB} \quad (7)$$

282 which can be seen plotted in Fig. 5 as a function of frequency.

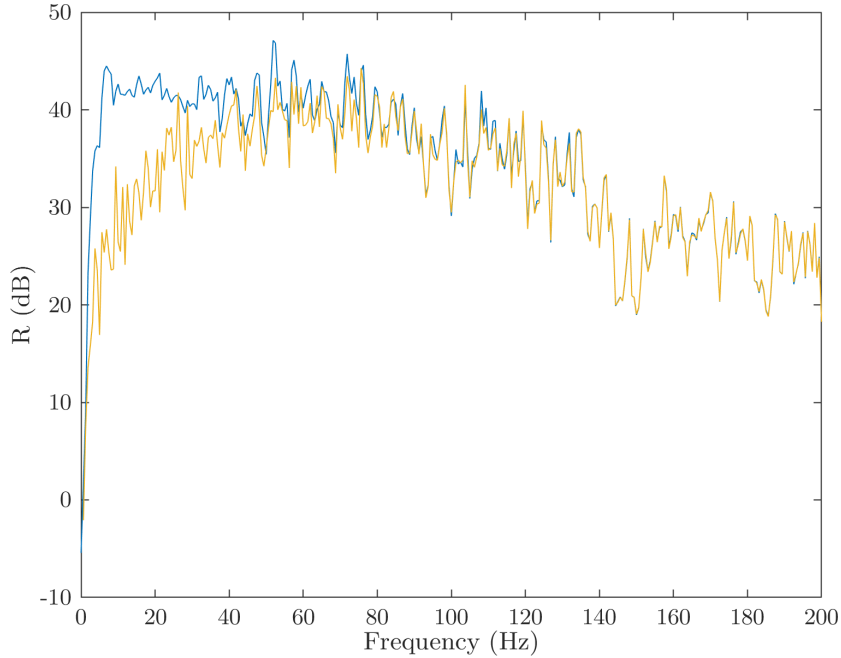


Figure 5: A plot of $R(f)$ obtained from $R(n)$ for both the established [11] and the improved frequency domain-based methods using five 1.6 s data lengths.

283 To obtain the data shown in Fig. 5, the reference channels for both tech-
 284 niques were processed identically and according to the technique presented
 285 in Fig. 4. These data show that the improved frequency domain-based tech-
 286 nique outperforms the established technique for frequencies below 100 Hz.
 287 However, for frequencies above 100 Hz, the difference is less noticeable and
 288 this is expected since the effect of the detrending is focused at lower frequen-
 289 cies. Quantitatively, this improvement translates to a seven times increase in
 290 R , with the established and the improved techniques obtaining $25.0^{+1.8}_{-1.3}$ dB
 291 and $33.5^{+1.2}_{-0.9}$ dB, respectively.

292
 293 While the correction performance was already substantial, this further
 294 improvement clearly shows the value of detrending after signal integration
 295 and is important for several reasons. Firstly, it yields a technique with much
 296 improved lower frequency performance which is likely to be beneficial for a
 297 range of important applications where low frequency, low-level vibrations are
 298 likely. Secondly, it resolves an important discrepancy previously observed
 299 when comparing frequency domain processing with a novel time domain-

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 300 based approach [16]. It should be noted that while previous work showed
 301 an eight times difference in the performance, the seven times difference in
 302 performance is not contradictory as the error reduction is dependant on the
 303 relative levels of the target and instrument vibration, therefore, is not consis-
 304 tent across setups. With performance inconsistency now resolved, the focus
 305 of the remainder of this paper will be a detailed assessment of the perfor-
 306 mance of the alternative time domain-based approach.
 307

308 **4. Theoretical generalisation of the synchronisation error on the** 309 **performance**

310 This section will work to establish an analytical model which relates the
 311 synchronisation error, $\Delta\tau$, to the quality of the correction. There will be
 312 a focus on time domain-based processing since the quality of the temporal
 313 alignment is limited to integer multiples of the time step. While interpola-
 314 tion could be used to upsample time domain data and enable sub-time step
 315 alignment, this is not always desirable. Therefore, this model can be used
 316 to predict the specifications of a system required to obtain high-quality time
 317 domain results based on the sampling frequency and the measured time delay.
 318

319 *4.1. Relating the error reduction to the synchronisation error*

320 To relate the synchronisation error, $\Delta\tau$, to the error reduction, R , it
 321 is assumed that $\Delta\tau$ is the primary factor which affects the quality of the
 322 corrected velocity estimate, MSE_{corr} . While other factors may also affect
 323 MSE_{corr} , this model is not concerned with them. Therefore, a relationship
 324 between MSE_{corr} and $\Delta\tau$ is required to relate $\Delta\tau$ to R .
 325

326 To do so, the corrected LDV signal, $v_{\text{corr}}(t)$, can be written as follows:

$$v_{\text{corr}}(t) = v_{\text{meas}}(t) - v_{\text{acc}}(t) \quad (8a)$$

327 where $v_{\text{meas}}(t)$ is the target velocity measured by the LDV and v_{acc} is the
 328 velocity of the LDV instrument itself, measured by the correction accelerom-
 329 eter. Rewriting Eq. (8a) to encapsulate the synchronisation error expressed
 330 as $v'_{\text{corr}}(t)$:

$$v'_{\text{corr}}(t) = v_{\text{meas}}(t) - v_{\text{acc}}(t + \Delta\tau) \quad (8b)$$

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331 The velocity error, $\Delta v_{\text{corr}}(t)$, can then be defined as the difference between
332 $v_{\text{corr}}(t)$ and $v'_{\text{corr}}(t)$:

$$\Delta v_{\text{corr}}(t) = v_{\text{acc}}(t) - v_{\text{acc}}(t + \Delta\tau) \quad (9)$$

333 Now a discrete Fourier expansion can be applied and, since $\Delta\tau$ is small,
334 the small angle approximation can also be applied:

$$v_{\text{acc}}(t) = \frac{a_0}{2} + \sum_{n=1}^N A_n \cos(n\omega_0 t) + B_n \sin(n\omega_0 t) \quad (10a)$$

$$\begin{aligned} v_{\text{acc}}(t + \Delta\tau) = \frac{a_0}{2} + \sum_{n=1}^N A_n \left(\cos(n\omega_0 t) - n\omega_0 \Delta\tau \sin(n\omega_0 t) \right) \\ + B_n \left(\sin(n\omega_0 t) + n\omega_0 \Delta\tau \cos(n\omega_0 t) \right) \end{aligned} \quad (10b)$$

335 where A_n and B_n are constants for each spectral line, a_0 is the DC com-
336 ponent and ω_0 is the spectral resolution. Substituting these expansions back
337 into Eq. (9) and simplifying the expression gives:

$$\Delta v_{\text{corr}}(t) = \Delta\tau \sum_{n=1}^N -A_n n\omega_0 \sin(n\omega_0 t) + B_n n\omega_0 \cos(n\omega_0 t) \quad (11)$$

338 Therefore, MSE_{corr} , or $\overline{\Delta v_{\text{corr}}(t)^2}$, can be expressed as:

$$\text{MSE}_{\text{corr}} = \frac{\Delta\tau^2}{2} \sum_{n=1}^N (A_n n\omega_0)^2 + (B_n n\omega_0)^2 \quad (12)$$

339 Inevitable sources of error other than synchronisation error mean that
340 the MSE will never be zero in practice. To account for this, an additional
341 term, c , is introduced:

$$\text{MSE}_{\text{corr}} = c + \frac{\Delta\tau^2}{2} \sum_{n=1}^N (A_n n\omega_0)^2 + (B_n n\omega_0)^2 \quad (13)$$

342 where c is a constant representing the lowest practically obtainable MSE
343 with a given setup. In order to relate this to the mean error reduction, Eq. (2)
344 can be rearranged into the following form:

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$$\frac{\text{MSE}_{\text{corr}}}{\text{MSE}_{\text{meas}}} = 10^{-\frac{R}{10}} = r \quad (14)$$

where r has been defined for convenience. Combining this with Eq. (13), the following can be written:

$$r(\Delta\tau) = \frac{c + \frac{\Delta\tau^2}{2} \sum_{n=1}^N (A_n n \omega_0)^2 + (B_n n \omega_0)^2}{\text{MSE}_{\text{meas}}} \quad (15)$$

The value of $r(\Delta\tau = 0)$ can then be described as “optimal”, and denoted by r_{opt} . Therefore, Eq. (15) can also be written as:

$$r(\Delta\tau) = r_{\text{opt}} + K\Delta\tau^2 \quad (16)$$

where r_{opt} and K have now incorporated all remaining constants. Both r_{opt} and K can be experimentally determined for a given setup.

Finally, an expression for $R(\Delta\tau)$ can be written by substituting Eq. (16) into Eq. (2):

$$R(\Delta\tau) = -10 \log_{10}(r_{\text{opt}} + K\Delta\tau^2) \quad (17)$$

However, to make use of this relationship, the synchronisation error must be derived and this differs for the frequency and time domain.

4.2. Frequency domain synchronisation error

The synchronisation error in the frequency domain, $\Delta\tau_f$, is simple since it only depends on how accurately the signal delay is known. Mathematically this can be defined as:

$$\Delta\tau_f = \tau_{\text{meas}} - \tau_{\text{true}} \quad (18)$$

where τ_{meas} is the measured time delay and τ_{true} is the theoretical true time delay. Practically, τ_{true} is the theoretical unknowable exact true time delay between the signals and τ_{meas} is determined using the procedure outlined in Section 3.1. Substituting this into Eq. (17) gives the following:

$$R(\tau_{\text{meas}}) = -10 \log_{10} \left(r_{\text{opt}} + K(\tau_{\text{meas}} - \tau_{\text{true}})^2 \right) \quad (19)$$

Therefore, this frequency domain model predicts there will be no sampling frequency dependence of the performance.

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369 *4.3. Time domain synchronisation error*

10 The synchronisation error in the time domain, $\Delta\tau_t(dt)$, is not only de-
11 pendent on the accuracy of the measured time delay, but also on the time
12 step. This is given by:
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$$15 \Delta\tau_t(dt) = \left\lceil \frac{\tau_{\text{meas}}}{dt} \right\rceil dt - \tau_{\text{true}} \quad (20a)$$

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19 where dt is the time step and $\lceil \dots \rceil$ denotes the nearest integer. Eq. (20a)
20 can now be written in terms of the sampling frequency, f_s , instead of the
21 time step since that is the adjustable acquisition parameter:
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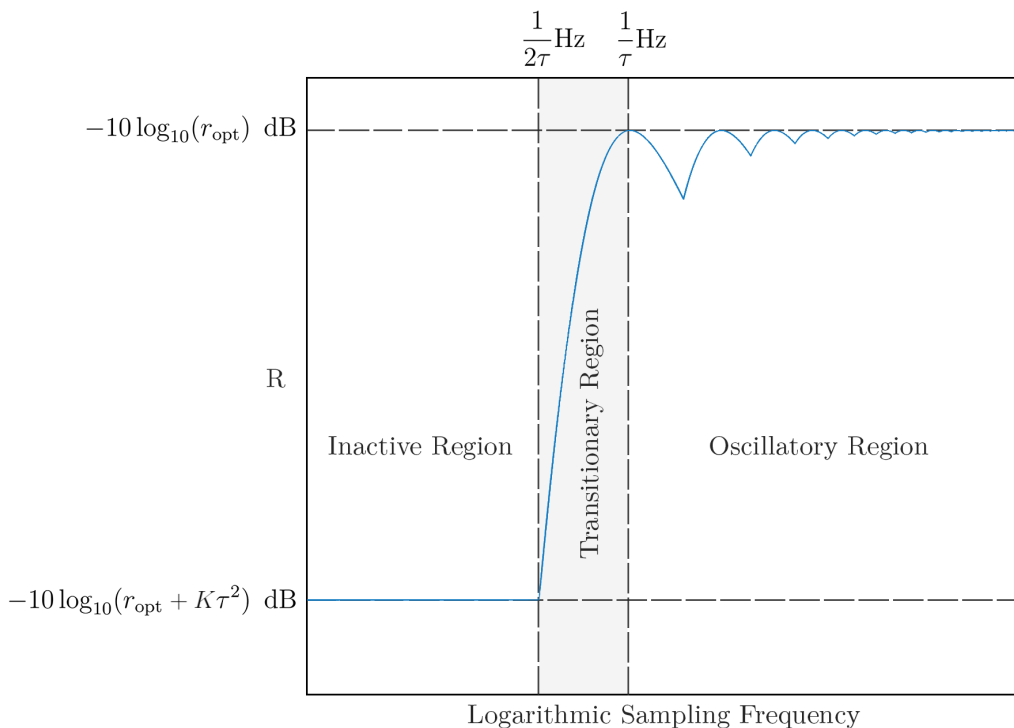
$$24 \Delta\tau_t(f_s) = \left\lceil \tau_{\text{meas}} f_s \right\rceil \frac{1}{f_s} - \tau_{\text{true}} \quad (20b)$$

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28 Since τ_{true} is the theoretically true value, it bares little practical signifi-
29 cance. Moving forward, it will be assumed that $\tau_{\text{meas}} \approx \tau_{\text{true}}$ so that the effect
30 of the sampling frequency alone on the error reduction can be thoroughly as-
31 sessed; both will now be denoted as τ . This assumption also results in both
32 positive and negative synchronisation errors having an equivalent negative
33 effect on the velocity estimate. Combining Eq. (17) and Eq. (20b) gives:
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$$37 R(f_s) = -10 \log_{10} \left(r_{\text{opt}} + K \left(\left\lceil \tau f_s \right\rceil \frac{1}{f_s} - \tau \right)^2 \right) \quad (21)$$

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41 The general form of the time domain model, with the significant features
42 labelled, can be seen in Fig. 6. As can be seen therein, there are three dis-
43 tinct regions. The first, “Inactive Region”, displays no sensitivity to the
44 sampling frequency. In this region the time step, dt , is too small for any
45 temporal alignment to take place; as such, no performance change occurs.
46 Temporal alignment becomes possible once the time step is smaller than $\frac{1}{2\tau}$,
47 representing the beginning of the “Transitional Region”. This region is
48 characterised by a large increase in the performance, as the decreasing time
49 step allows for increasingly more accurate temporal alignment. The third
50 and final region, the “Oscillatory Region”, is characterised by oscillations
51 in performance which decrease in amplitude as the frequency increases. The
52 peaks of these occur at integer multiples of τ^{-1} Hz, since these locations are
53 where τ becomes divisible by an integer number of time steps, leading to an
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392 increase in the accuracy of the temporal alignment. Similarly, the perfor-
393 mance troughs occur at integer multiples of $0.5\tau^{-1}$ Hz
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42 Figure 6: A general plot of the time domain error reduction model as a function of the
43 sampling frequency with three distinct regions labelled. Higher values on the vertical scale
44 represent better performance.

45
46 Since a continuous range of sampling frequencies is rarely available, a
47 more practically relevant example is Eq. (21) plotted at the sampling frequen-
48 cies available on the Simcenter SCADAS Mobile data acquisition system, as
49 seen in Fig. 7. While the highest sampling frequency, 204.8 kHz, shows a sig-
50 nificant reduction in error, the same performance could have been achieved
51 by using 8.192 kHz, 16.384 kHz or 40.960 kHz. In fact, the aforementioned
52 frequencies have a error reduction 0.03 dB higher than the highest sampling
53 frequency. This shows that, even without measuring values for the constants
54 K and r_{opt} , the time domain model can still predict the optimal sampling
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404 frequency to maximise performance.

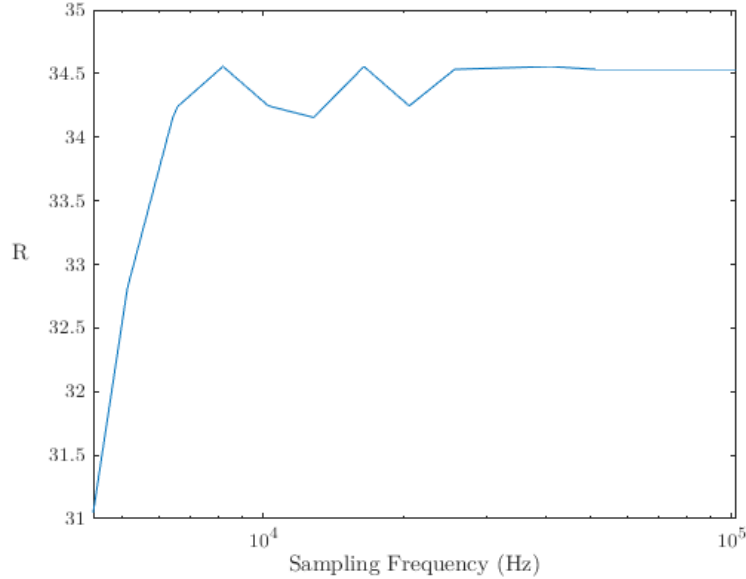


Figure 7: A plot of error reduction as a function of the sampling frequencies available on the Simcenter SCADAS Mobile data acquisition system. This is plotted with values of $K = 35 \times 10^3 \text{ s}^{-2}$ and $r_{\text{opt}} = 350 \times 10^{-6}$.

405 4.4. Time domain constants determination

406 Since the main use of the time domain model is to enable the informed se-
 407 lection of the sampling frequency, knowledge of the constants is not necessary.
 408 However, in order to validate the time domain model it will be compared to
 409 an experimentally measured $R(f_s)$, denoted by $R_{\text{ex}}(f_s)$, meaning the con-
 410 stants are required since they affect the model's relative proportions in the
 411 vertical axis.

412 The first constant, r_{opt} , is calculated using:
 413

$$r_{\text{opt}} = 10^{-\frac{\max(R_{\text{ex}})}{10}} \quad (22)$$

414 where $\max(R_{\text{ex}})$ is the maximum value of $R_{\text{ex}}(f)$. Similarly, K , can be
 415 calculated using:

$$K = \frac{r_{\text{opt}} - 10^{-\frac{R_{\text{ex}}(f_s)}{10}}}{\left(\left[\tau_{\text{meas}} f_s \right] \frac{1}{f_s} - \tau_{\text{meas}} \right)^2} \quad \text{for} \quad f_s > \frac{1}{2\tau_{\text{meas}}} \text{ Hz} \quad (23)$$

where all symbols are as previously defined. The sampling frequency here must be larger than $\frac{1}{2\tau_{\text{meas}}}$ Hz as the time domain model does not predict any behaviour in the Inactive Region so scaling using these data will lead to erroneous predictions.

5. Experimental validation of the time domain model

This section has two aims: firstly, to show that the improvements made to the established frequency domain-based technique close the performance gap when compared to the time domain-based technique; secondly, to show their relative performances over a range of frequencies and, by doing so, confirm the time domain model for $R(f_s)$, given by Eq. (21). This will provide the user with two instrument vibration correction techniques when faced with stationary and transient signal types. The experimental arrangement used in the following was common with that used in Section 2.

5.1. Data collection and processing

To validate the time domain model presented in Eq. (23), the error reduction needs to be characterised against the sampling frequency and thus data is acquired at the highest available sampling frequency of 204.8 kHz and iteratively downsampled to simulate acquisition at lower sampling frequencies. The process was implemented in MATLAB and Fig. 8 illustrates this code for the time domain and improved frequency domain-based techniques.

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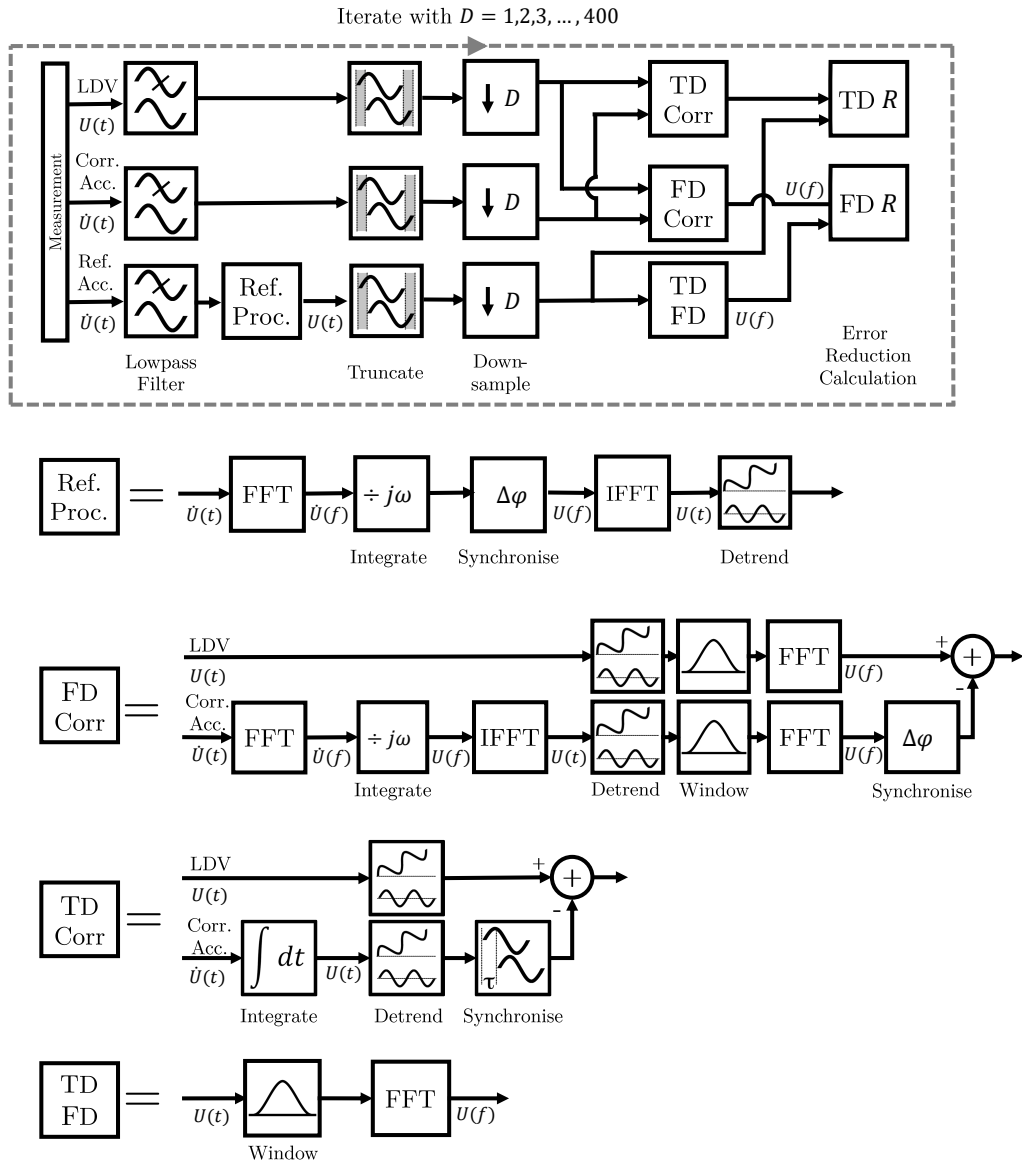


Figure 8: A schematic of the code used to characterise the error reduction as a function of sampling frequency. Where D is a downsampling factor and “TD Corr” and “FD Corr” represent both the time domain [16] and improved frequency domain-based correction techniques, respectively.

439 As shown in Fig. 8, measured data are immediately low-pass filtered using
 440 a finite impulse response digital lowpass filter with a 200 Hz cut-off frequency
 441 which coincides with the maximum frequency of the vibration. Any spuri-

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442 ous higher frequency signal content which might otherwise have been aliased
443 into the frequency range of interest following the downsampling is thereby
444 rejected. Following this, the reference accelerometer signal is then subjected
445 to the same frequency domain detrending and time synchronisation steps as
446 that previously described. This special treatment is present purely to ensure
447 that the reference signal is as close to a ‘true’ signal as possible, in terms of
448 both the integration accuracy and synchronisation error.

449
450 Since the Fourier Transform implicitly assumes that a signal is periodic,
451 the phase shift would have caused a portion of the signal at the beginning or
452 the end of the reference accelerometer signal to wrap around to the opposite
453 end of the signal, depending on which signal is lagging. To fix this, all sam-
454 ples in this ‘wrapped’ region are removed from all three transducer signals
455 in the stage named “Truncation” shown in Fig. 8. Following this, the sig-
456 nals are downsampled by taking each D th sample from the original signals,
457 simulating a lower sampling frequency acquisition. The penultimate stages
458 named “TD Corr” and “FD Corr” represent the two correction algorithms.
459 This process was looped in the code with $D = 1, 2, \dots, 400$, giving a mini-
460 mum sampling frequency of 512 Hz and a total of 400 data points for each
461 correction technique. The final output is, therefore, two data sets describing
462 the performance of each correction algorithm as a function of the sampling
463 frequency.

464 465 *5.2. Model validation and sample rate dependent performance assessment*

466 The frequency domain model, given by Eq. (19), predicts no sampling
467 frequency dependence as sub-time step synchronisation is possible in the fre-
468 quency domain. However, the time domain model, given by Eq (21), predicts
469 a reasonably strong dependence due to this time step synchronisation lim-
470 itation. Comparing Figures 9 and 6, the three previously defined regions
471 are clearly identifiable from the experimental data. In particular, the impor-
472 tant Transitional Region is clearly shown; this is where the performance of
473 the time domain-based technique “overtakes” that of the frequency domain-
474 based approach.

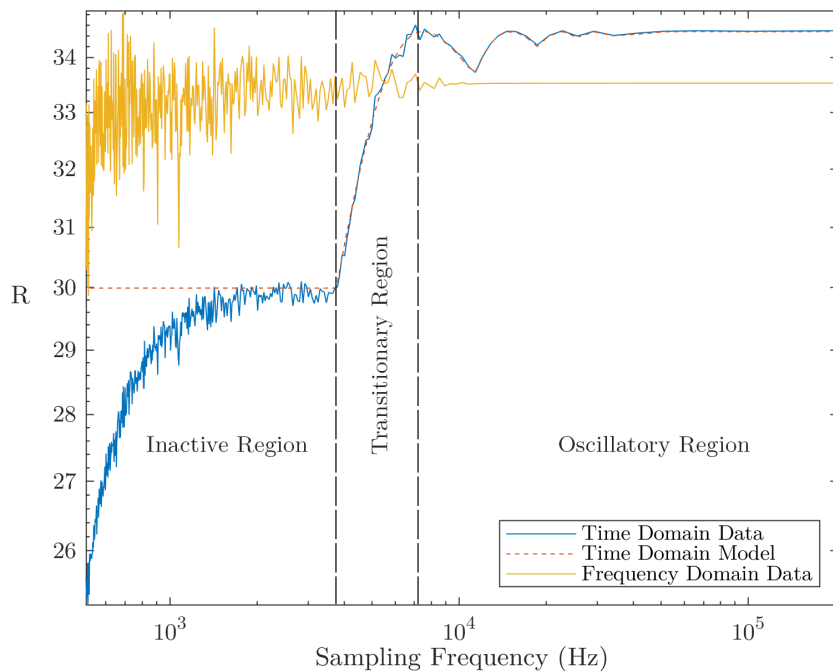


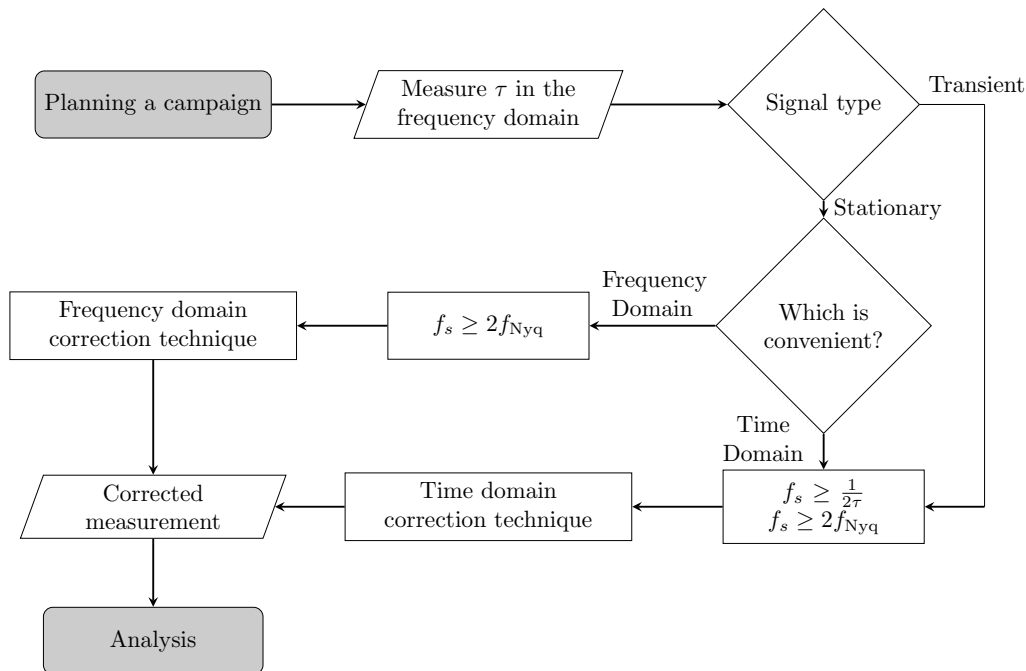
Figure 9: A plot of the experimental error reduction as a function of the sampling frequency for the improved frequency domain and the time domain-based techniques. The error reduction is calculated using Eq. (2) and the domain-specific formulations for the MSE. The time domain model, given by Eq. (21), is also plotted for validation purposes.

476 In order to validate the time domain model, the time domain data can
 477 be compared to the time domain model in Fig. 9; from this, two shortcom-
 478 ings can be seen. Firstly, the time domain model does not predict the smaller
 479 noise-like fluctuations visible at frequencies below 10 kHz in the time domain
 480 data. However, these fluctuations are exhibited by both correction techniques
 481 and are likely caused by different factors as they do not occur at common
 482 frequencies. There was no attempt to model these fluctuations, therefore,
 483 this is no major shortcoming of either model. Secondly, the time domain
 484 model fails to predict the behaviour in the Inactive Region. The three times
 485 increase seen in the time domain experimental data is likely caused by a de-
 486 crease in the quality of integration at the lower sampling frequencies when
 487 using the cumulative trapezoidal method. The time domain model does not
 488 capture this behaviour as it only considered the effects of temporal alignment.

489
 490 However, as can be seen in Fig. 9, the time domain model almost ex-

491 actly describes the experimental data in the Transitional and Oscillatory
 492 Regions, which are the regions of interest here. With the time domain model
 493 validated, its value is its ability to determine when the time domain technique
 494 will perform optimally, based on the time delay estimate and the sampling
 495 frequency. That is, significant improvements will be made to the quality of
 496 the correction if a data acquisition system is used with a sampling frequency
 497 larger than $\frac{1}{2\tau}$ Hz. Similarly, any sampling frequency larger than τ^{-1} Hz
 498 will not yield any substantial increase in performance and is therefore un-
 499 necessary. Also, if possible, a sampling frequency should be selected close to
 500 a performance peak located at $\frac{n}{2\tau}$ Hz with $n = 1, 2, 3, \dots$

501



47 Figure 10: A flow chart describing when to use either the time domain [16] or the improved
 48 frequency domain-based technique described herein. The outcome is based on the sampling
 49 frequencies available to the user and τ .

502 Comparing the performances of the two techniques in the region $f_s > \frac{1}{2\tau}$
 503 Hz, as shown in Fig. 9, there is about a decibel of performance gain to be
 504 made by using the time domain-based technique. The exact reason for the
 505 difference in performance is yet to be determined. However, since this is

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506 small, either technique could be used with minimal difference in achieved
507 performance. Similarly, the user could select a sampling frequency close to
508 a performance peak to optimise the time domain-based technique perfor-
509 mance. However, the benefit of this is also marginal. Fig. 10 summarises
510 these generalised findings based on the vibration signal type and the sam-
511 pling frequency, advising the user which is the most appropriate technique
512 to use for a given measurement campaign. When the vibration is stationary
513 in nature, the user can use either technique. When the vibration is transient
514 in nature, the user must use the time domain-based technique and select an
515 appropriate sampling frequency to optimise performance.
516

517 6. Conclusions

518 Recent advances in the application of LDVs to measurement campaigns
519 in which the instrument sensor head is itself subject to vibration have lead
520 to an increasing number of techniques for the correction of the measured
521 signals. Practical implementation of these techniques involves the determi-
522 nation of the sensor head vibration and subtraction of this in post-processing.
523 Extension from lab to field-based measurements has further necessitated the
524 conception and development of novel time domain-based processing tech-
525 niques for vibration signals that are transient in nature. Initial investigations
526 showed that, for common signals, these alternative techniques significantly
527 outperformed previously established frequency domain equivalents.

528 Firstly, therefore, this work aimed to close this previously observed per-
529 formance gap with an improved *frequency domain* based technique being
530 developed. A seven times performance increase was obtained by applying a
531 modified frequency domain-based technique which included a detrending step
532 prior to implementation of the correction processing. To make this detrend-
533 ing possible, it is necessary that no window is applied to the sampled data
534 until after detrending. Particular improvement was shown to be found in fre-
535 quencies below 100 Hz which is arguably a major benefit since applications
536 of interest for such techniques are expected to be focused in this frequency
537 range.

538 Since the required correction measurements are typically obtained using
539 accelerometers, in addition to the requisite integration, it is typical that sig-
540 nal synchronisation is necessary due to signal conditioning differences. It is
541 well known that the quality of the signal synchronisation will contribute to

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10 542 the correction performance. Therefore, when working in the *time domain*
11 543 and when interpolation is not desirable, the sampling frequency would then
12 544 contribute to the synchronisation error as time shifts are only possible in
13 545 units of the time step. As such, a model describing the relationship between
14 546 the synchronisation error and the performance was derived and formulated
15 547 in terms of the sampling frequency and the error reduction. To validate this
16 548 model, code was written to obtain the sampling frequency dependence of the
17 549 performance by iteratively downsampling high-sample rate data to simulate
18 550 acquired data at a range of sampling frequencies. The various correction
19 551 techniques were then tested on these data to experimentally obtain the rela-
20 552 tionship derived in the model with excellent agreement found.

23 553 Given two viable and equally effective correction techniques, each with
24 554 their own set of requirements, a framework was developed to allow the user to
25 555 conveniently select the appropriate correction technique, taking into account
26 556 the specifics of the vibration measurement of interest. Whichever the required
27 557 technique, the significant contribution of this work is to enable the user to
28 558 optimally specify the data acquisition parameters. This enables definition
29 559 of the optimal hardware characteristics required for a given measurement
30 560 campaign, important for efficient and practical integration of such sensor
31 561 solutions.

36 562 7. Acknowledgements

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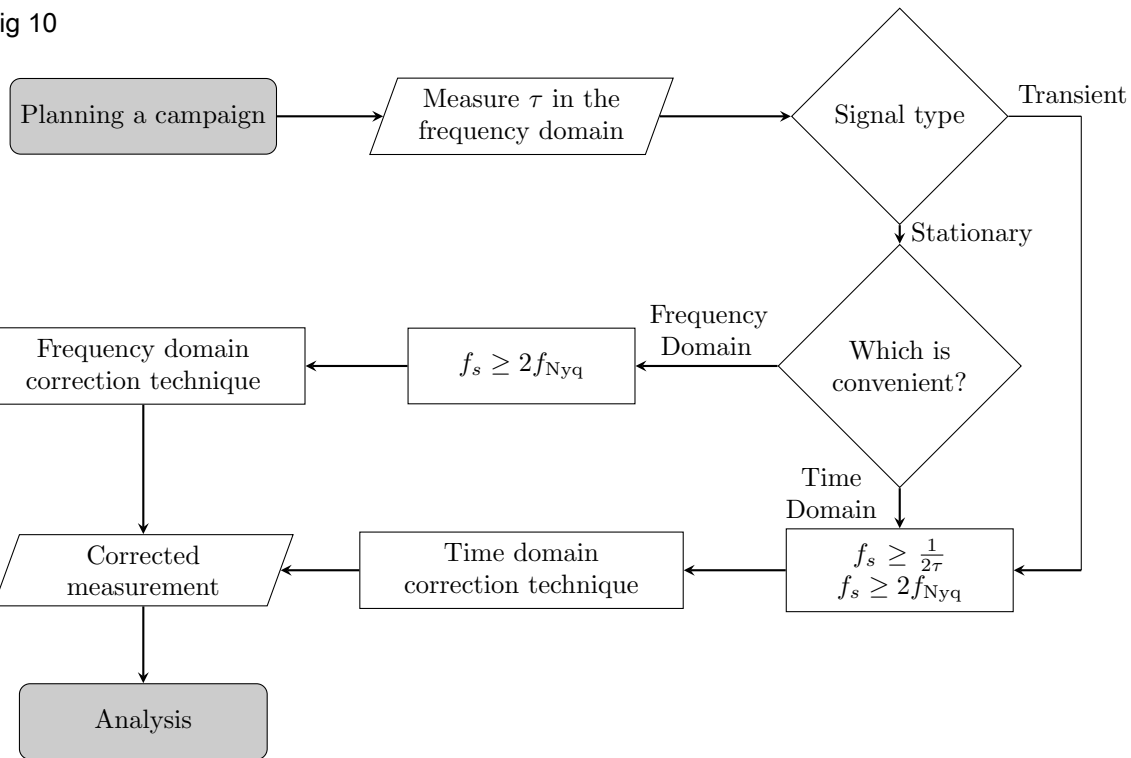


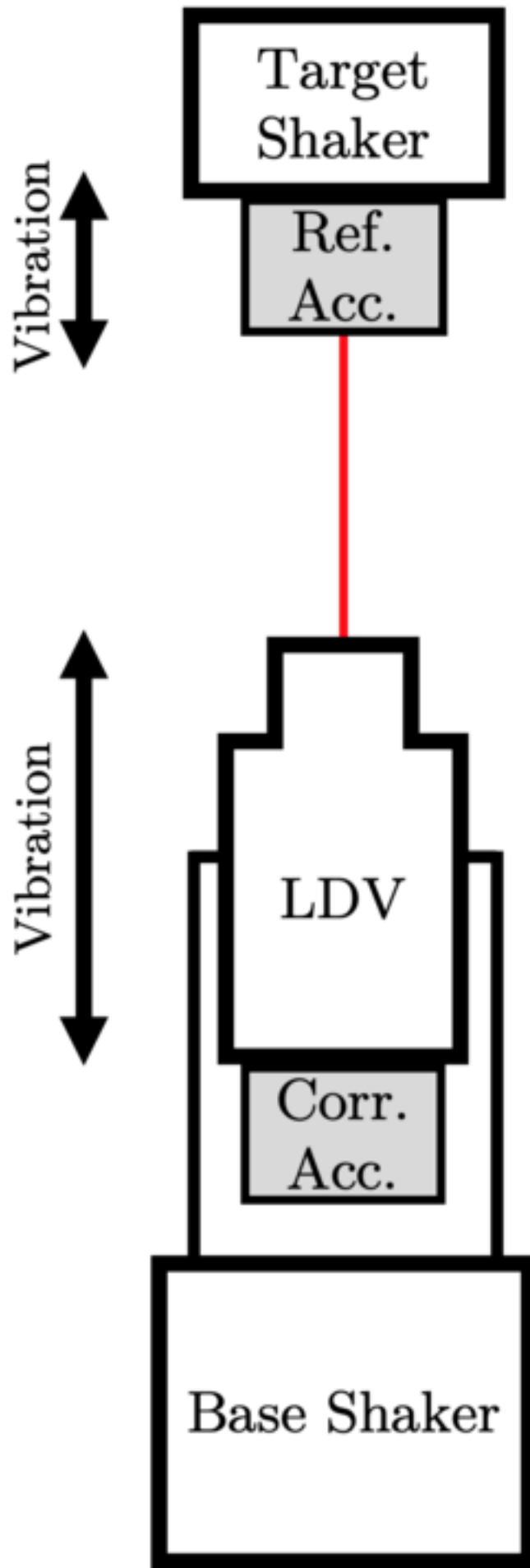


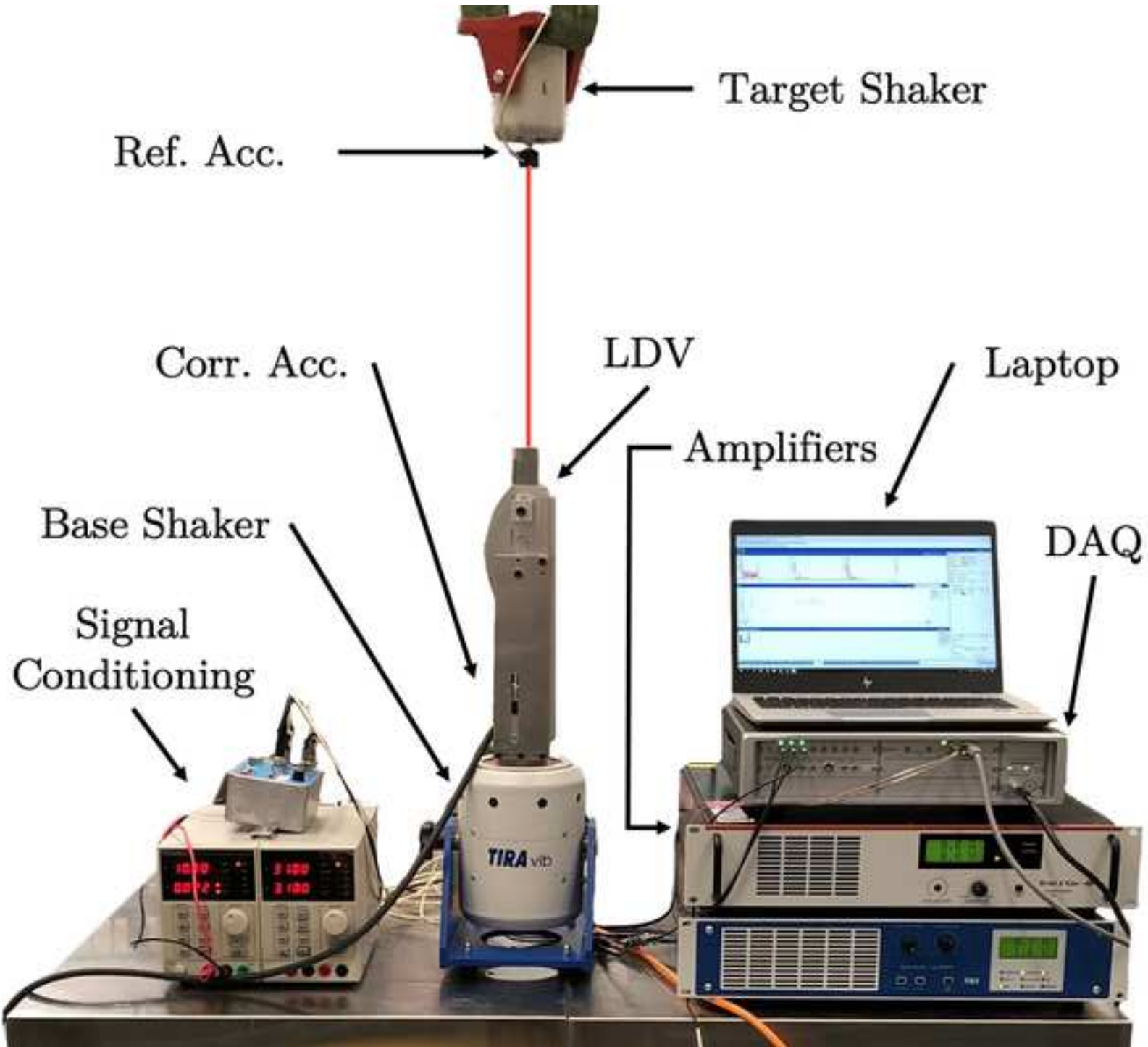
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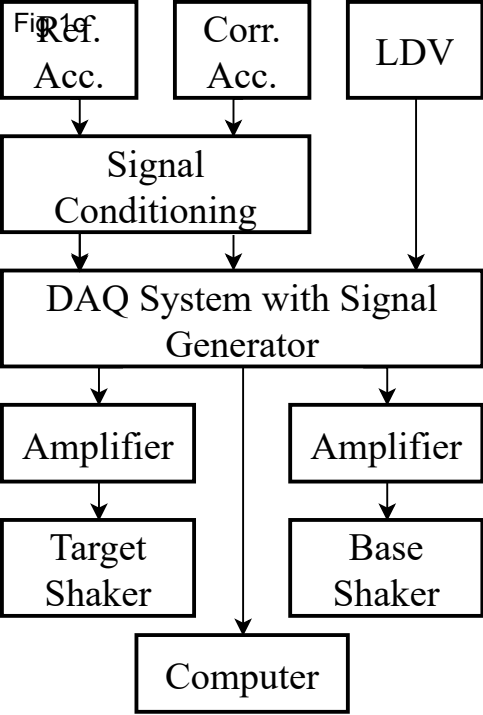


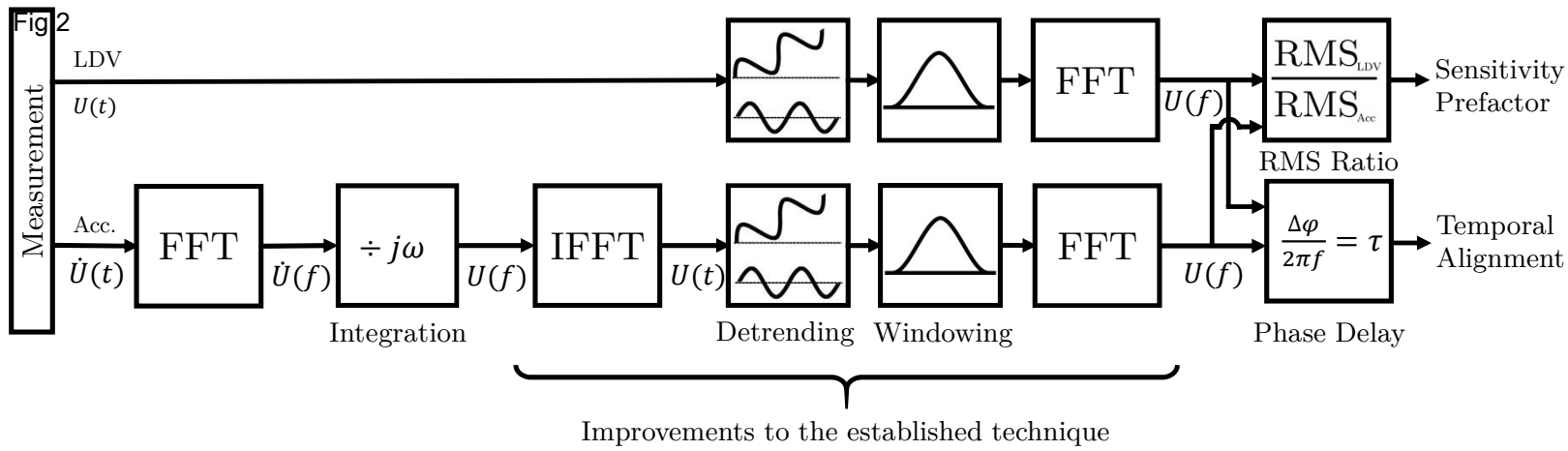
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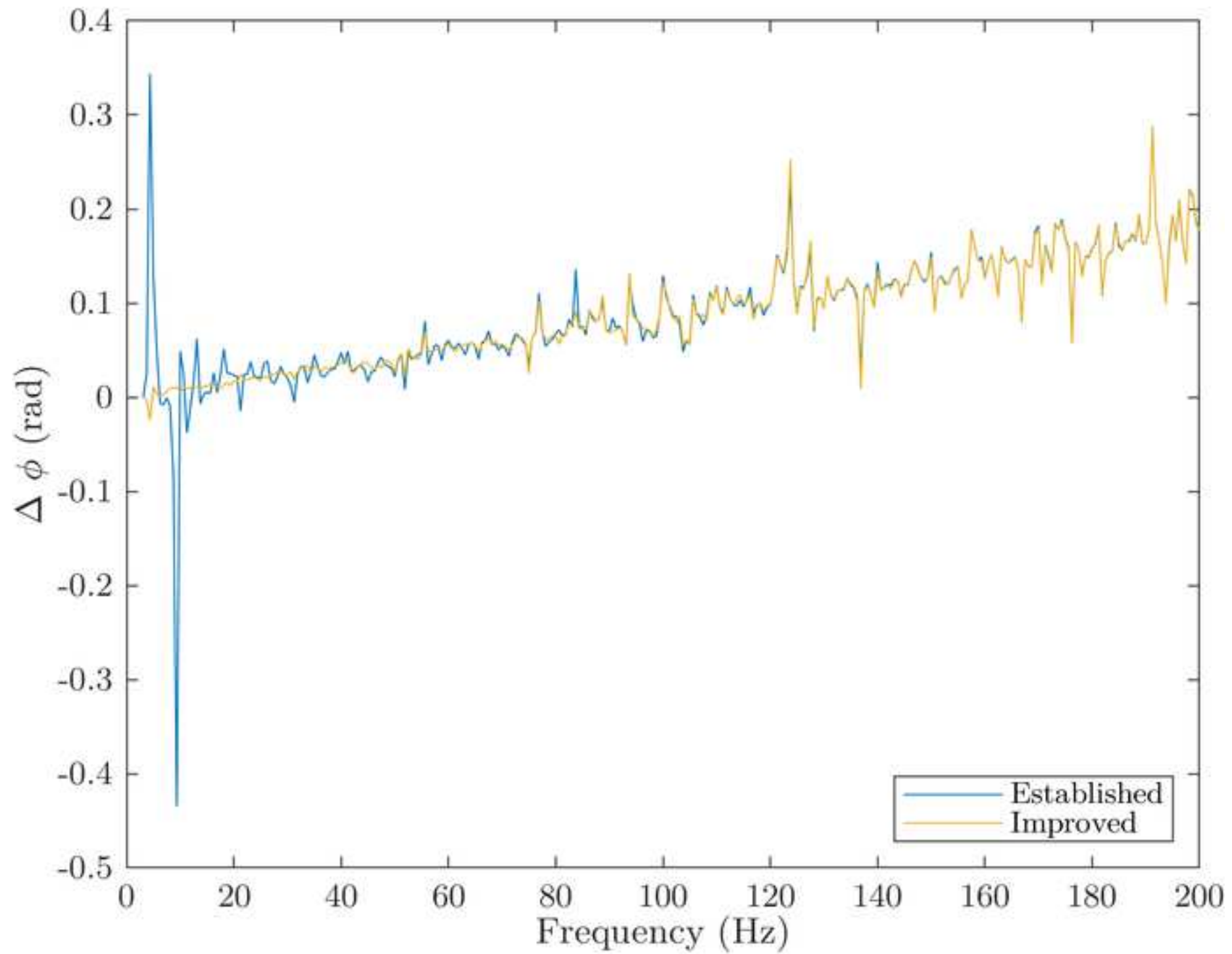


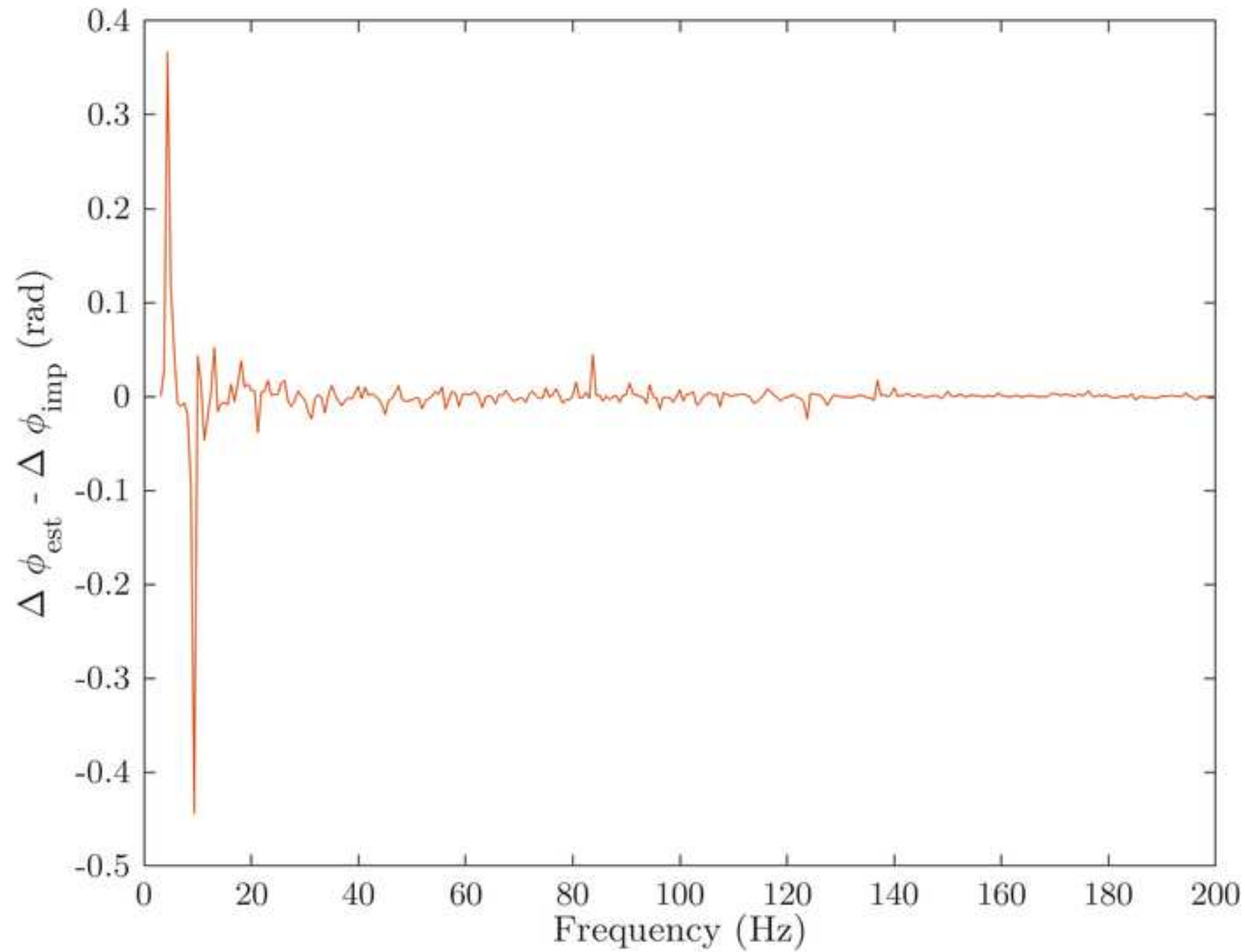












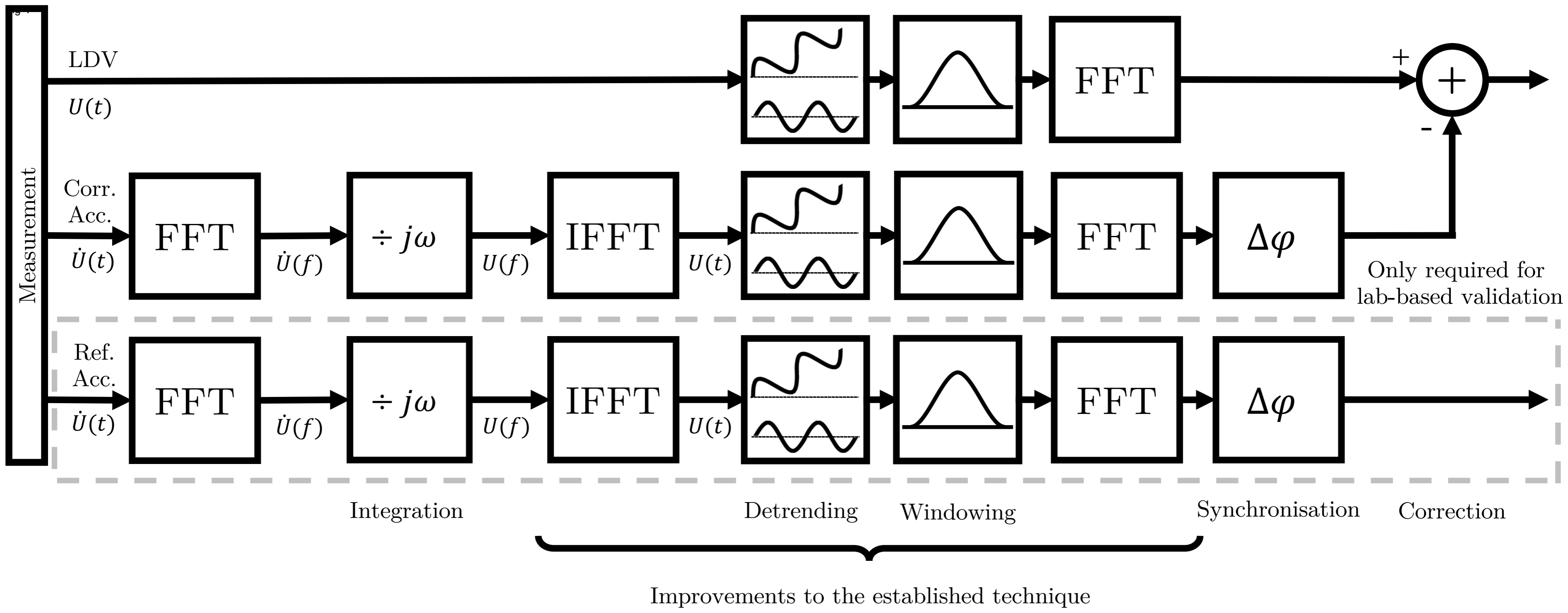
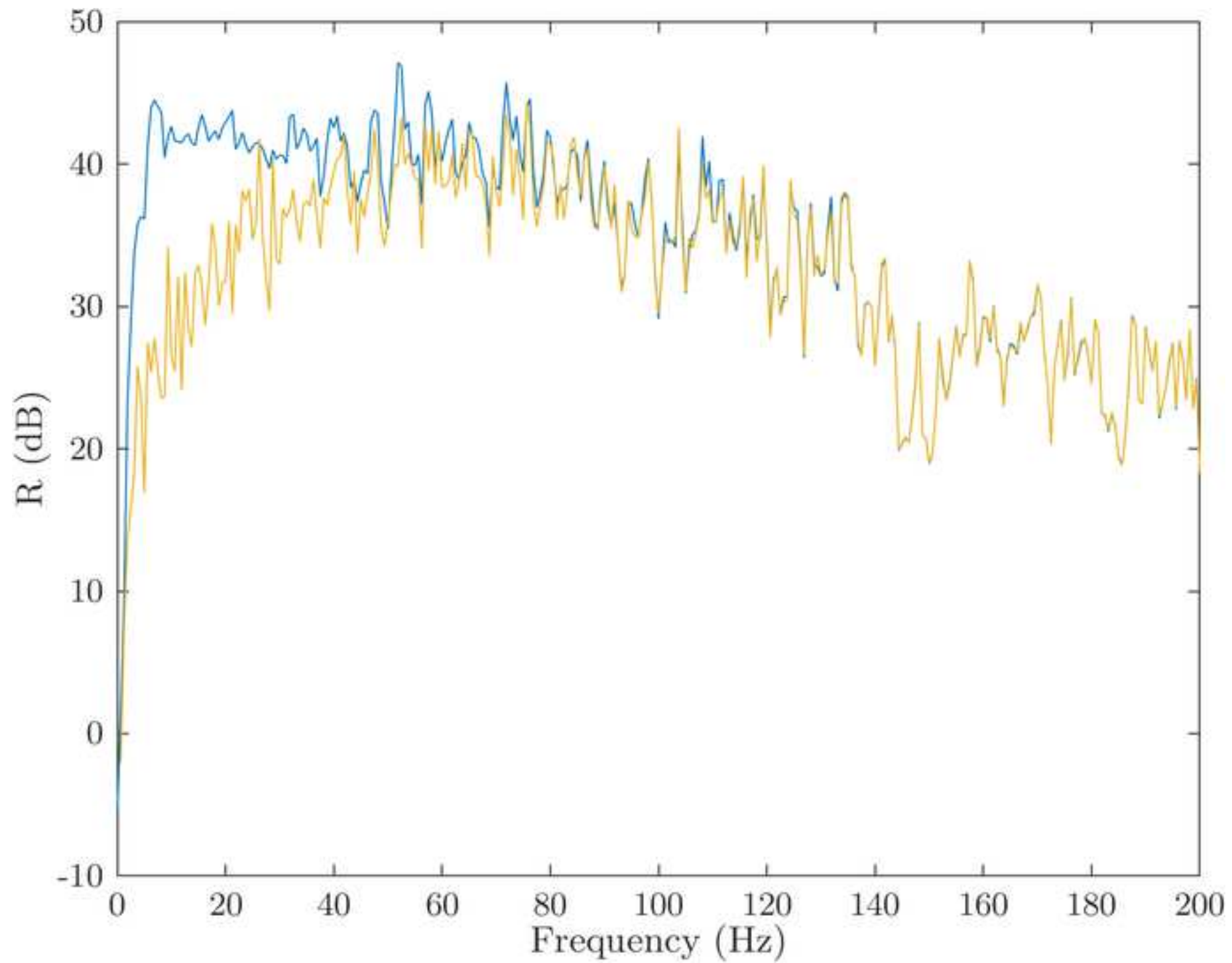
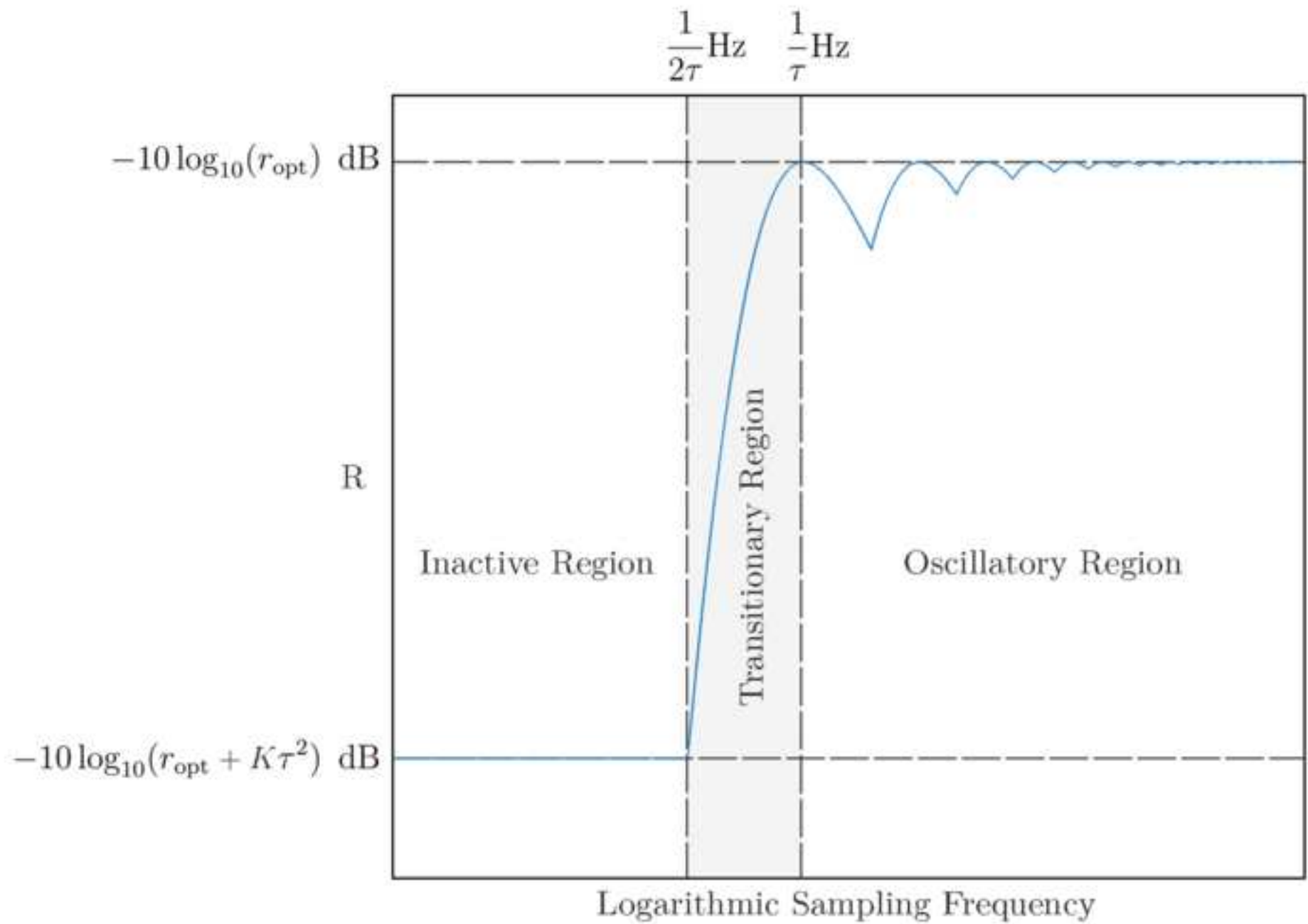
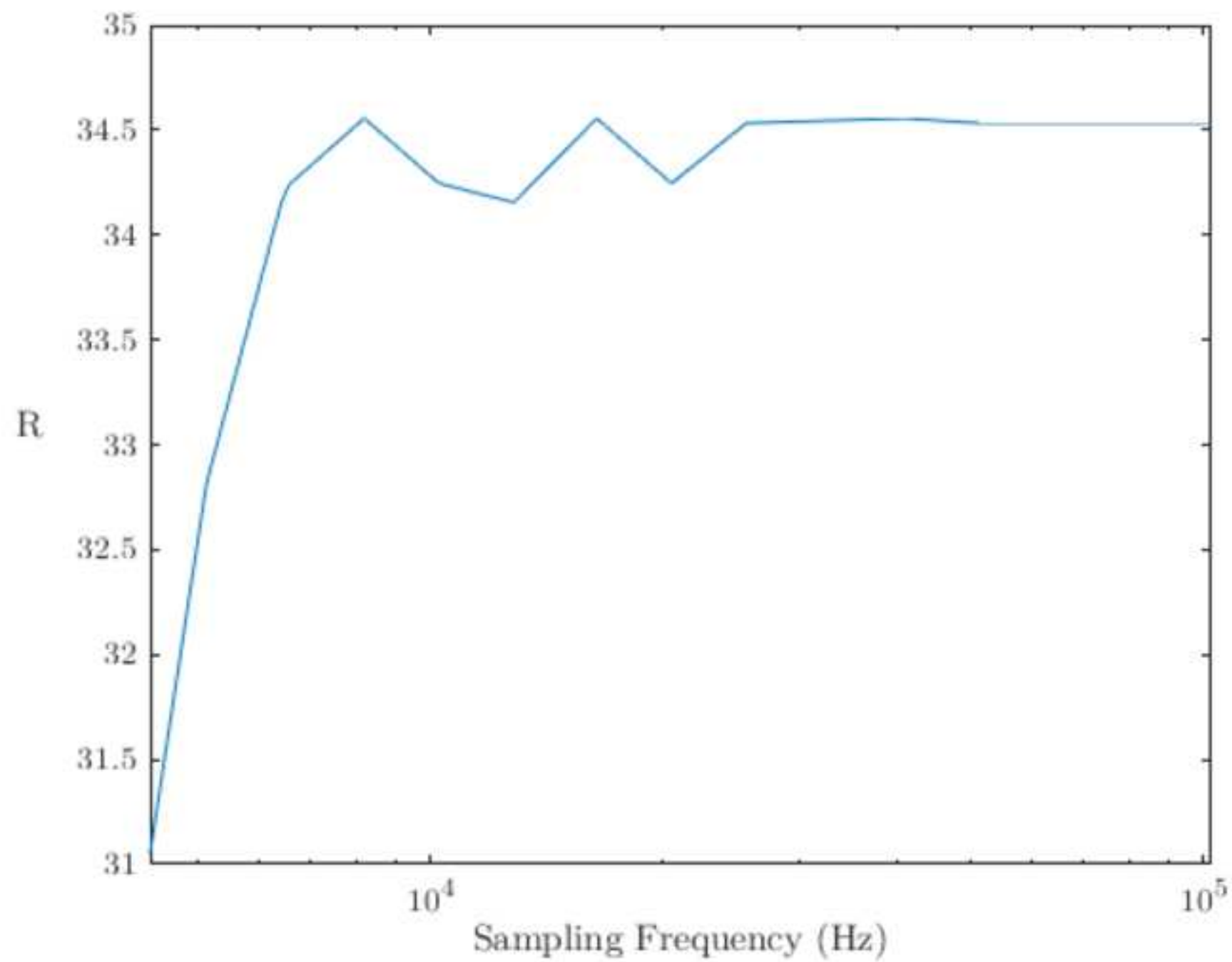


Fig 5







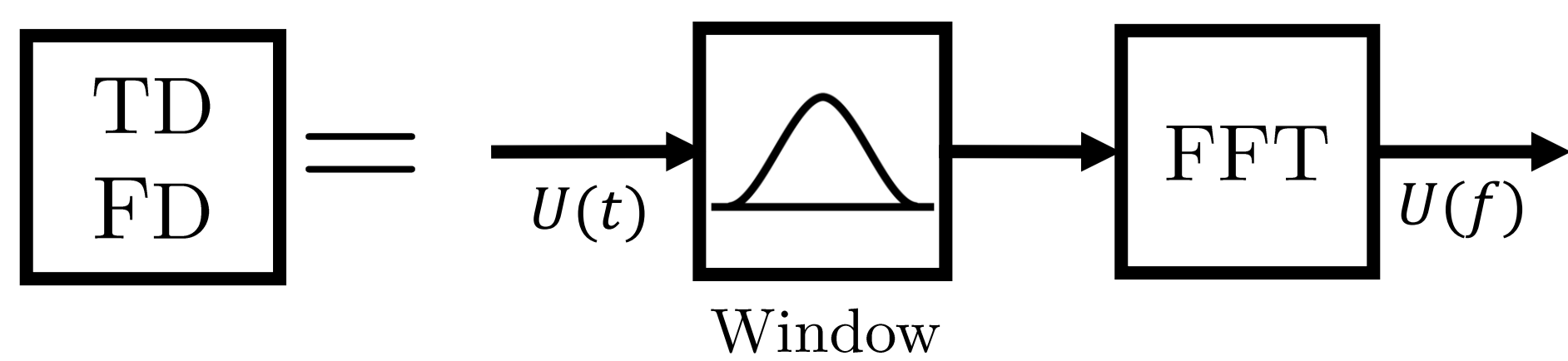
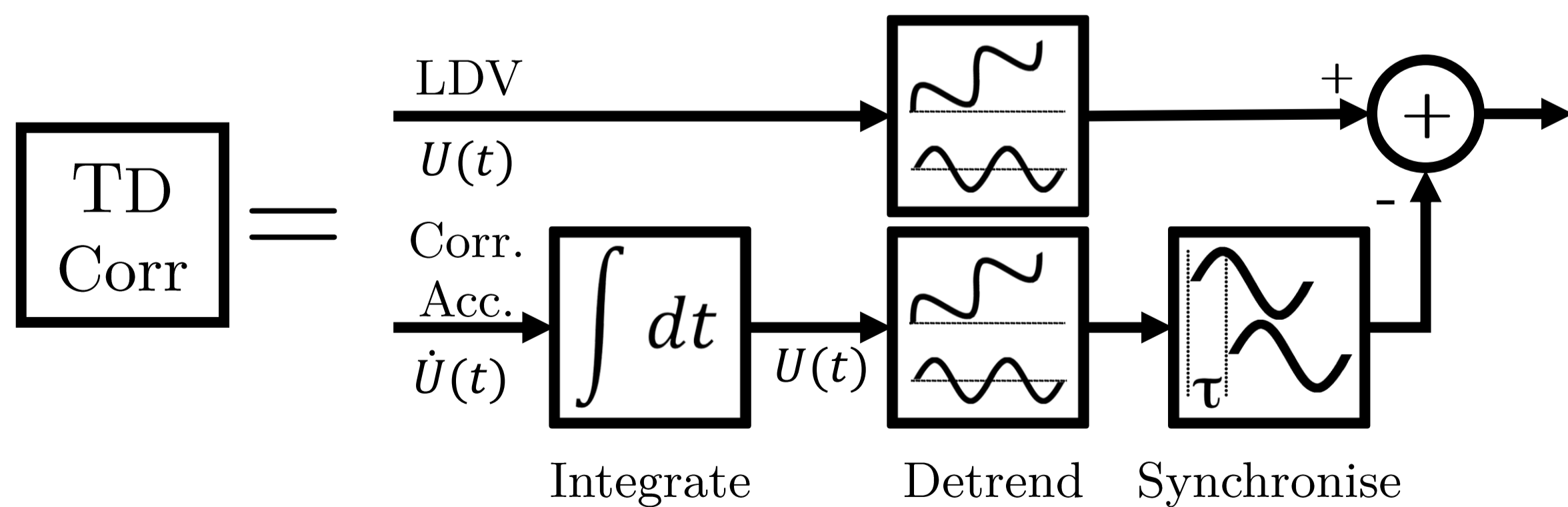
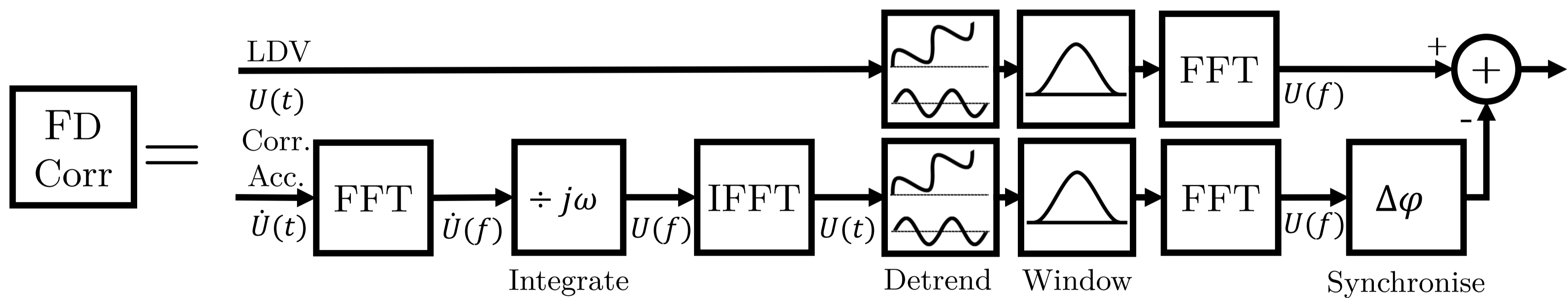
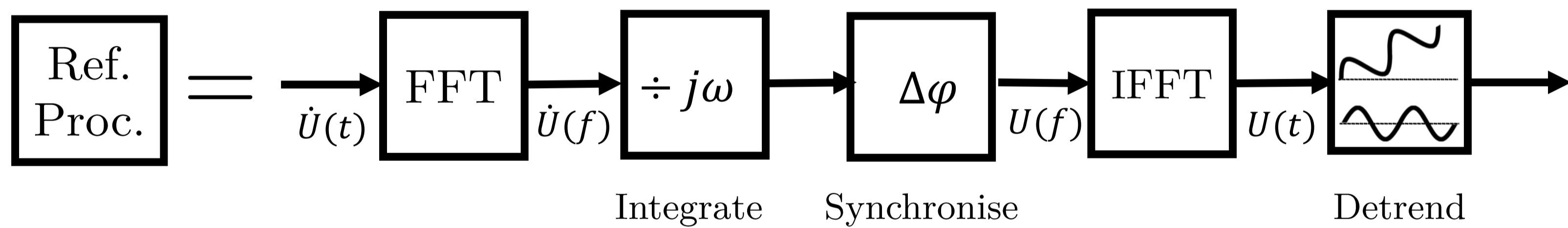
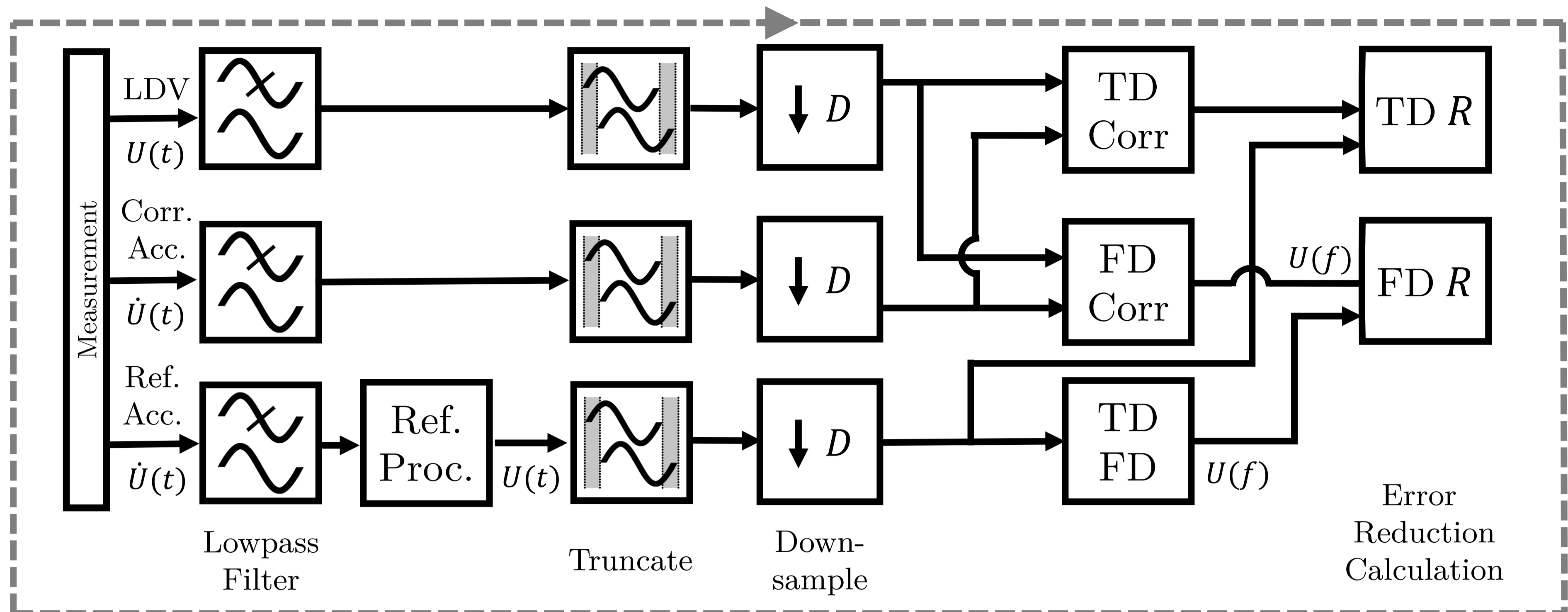
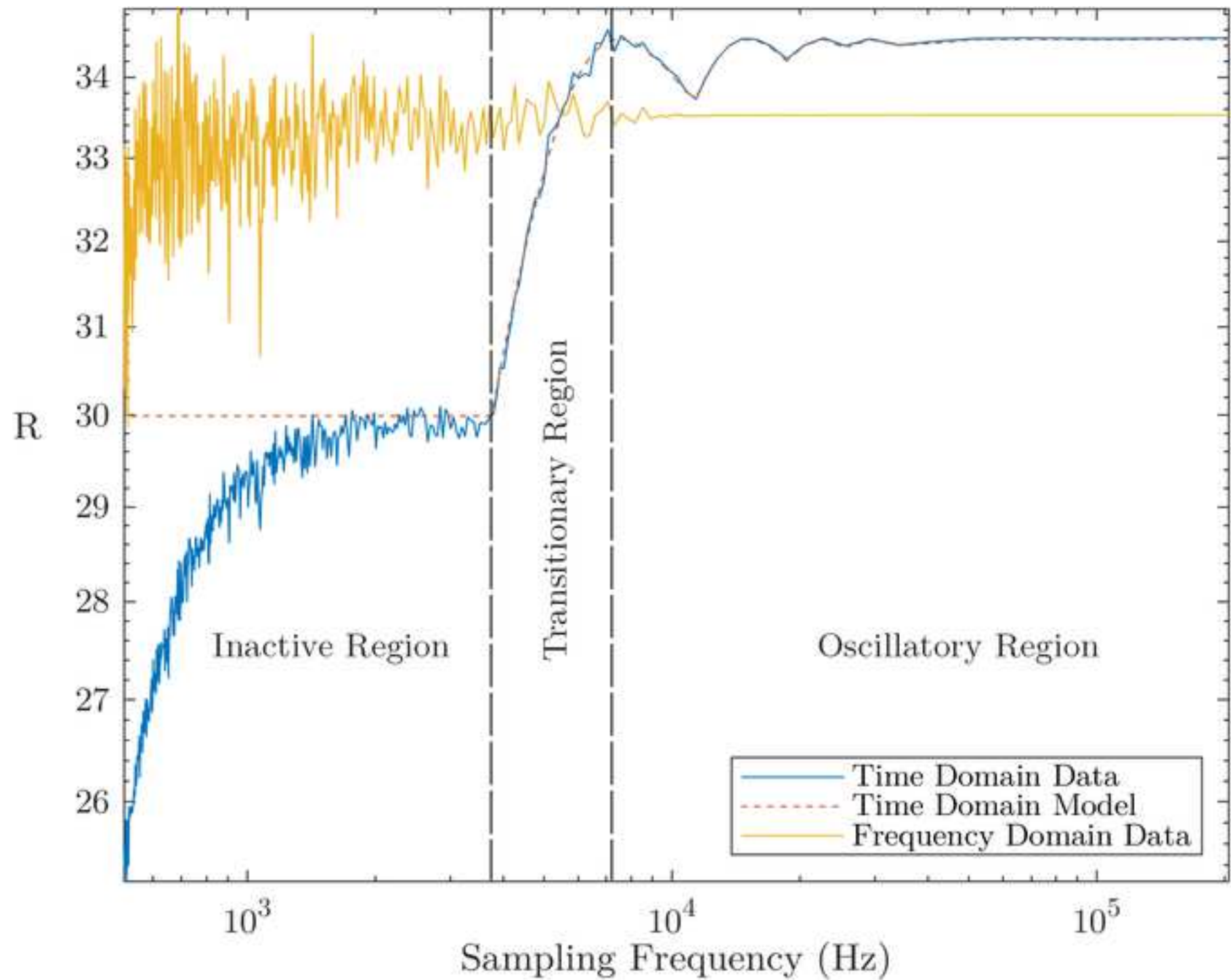
Iterate with $D = 1, 2, 3, \dots, 400$ 

Fig 9



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Abdel Darwish: Conceptualisation, Methodology, Software, Validation, Formal analysis, Investigation, Resources, Data Curation, Writing - Original Draft, Writing - Review & Editing, Visualisation and Project administration.

Benjamin Halkon: Conceptualisation, Methodology, Software, Validation, Formal analysis, Investigation, Resources, Data Curation, Writing - Original Draft, Writing - Review & Editing, Visualisation, Supervision, Project administration and Funding Acquisition.

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Sebastian Oberst: Methodology, Validation, Resources, Writing - Review & Editing, Supervision and Funding Acquisition.

Robert Fitch: Conceptualisation, Resources, Validation, Writing - Review & Editing, Supervision and Funding Acquisition.

Declaration of interests

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

The authors declare the following financial interests/personal relationships which may be considered as potential competing interests: